

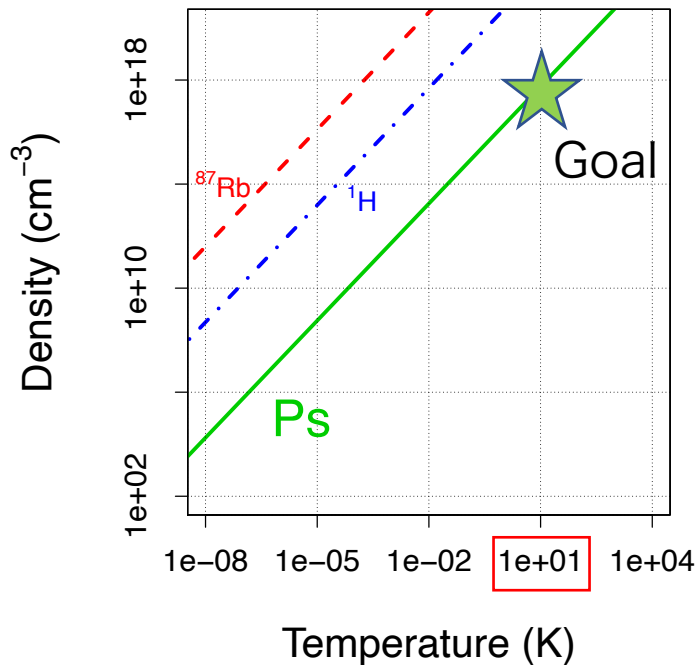
# Development of a High-Resolution Probing Laser Suited for Cold Positronium Spectroscopy

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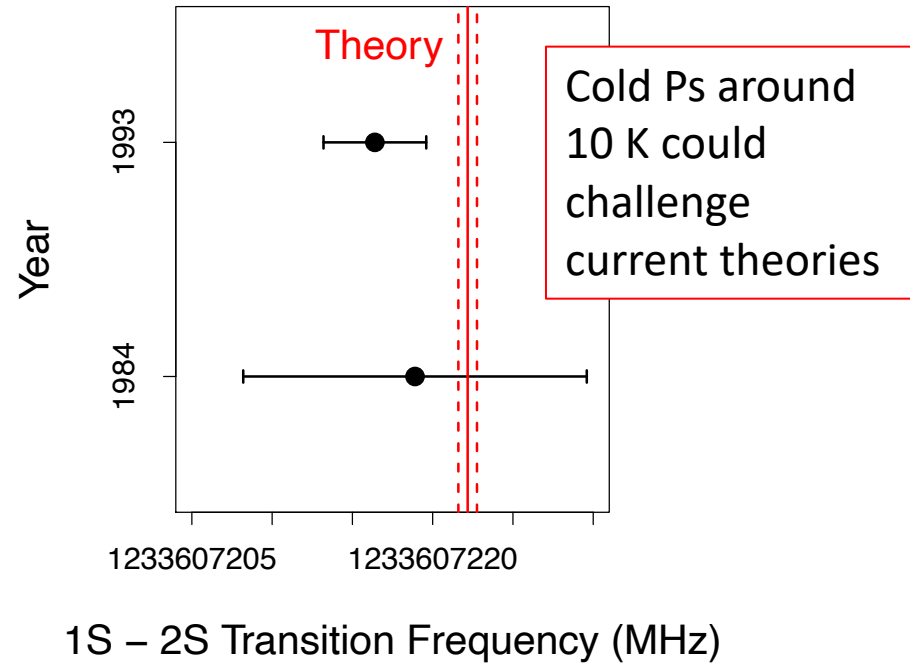


# Significance of Cold Ps

Realization of the first Bose-Einstein Condensation by an antimatter containing atom



Precision spectroscopy of transition frequencies



M. S. Fee *et al.*, *Physical Review A* **48**, 192 (1993).  
K. Danzmann, M. S. Fee, and S. Chu, *Physical Review A* **39**, 6072 (1989)  
A. Czarnecki, K. Melnikov, and A. Yelkhovsky, *Physical Review A* **59**, 4316 (1999).

Realizing cold Ps around 10 K is of great importance

# Temperature Measurement Methods of Cold Ps needs to be Established

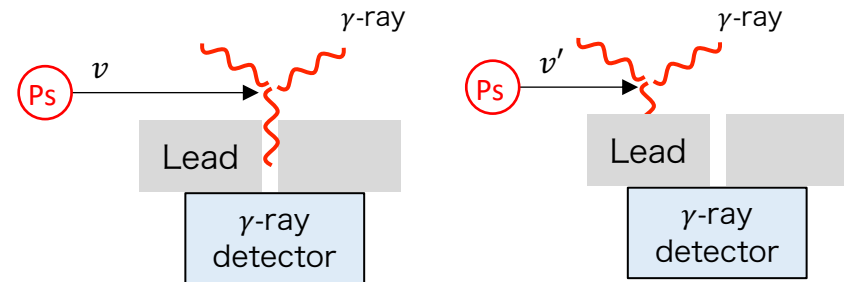
## Typical Temperature measurement methods

- Time-of-flight (TOF)

Ps mean velocity:

$1.4 \times 10^4$  m/s @10 K

$2.0 \times 10^4$  m/s @20 K



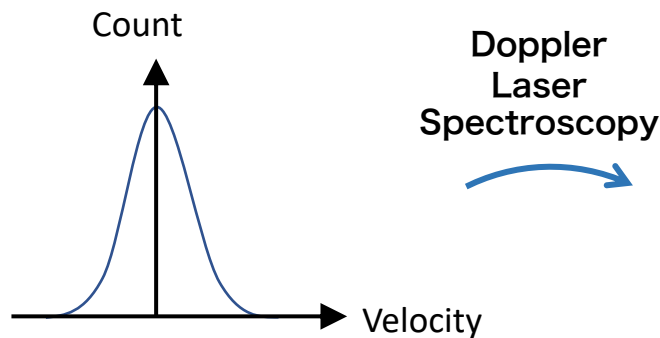
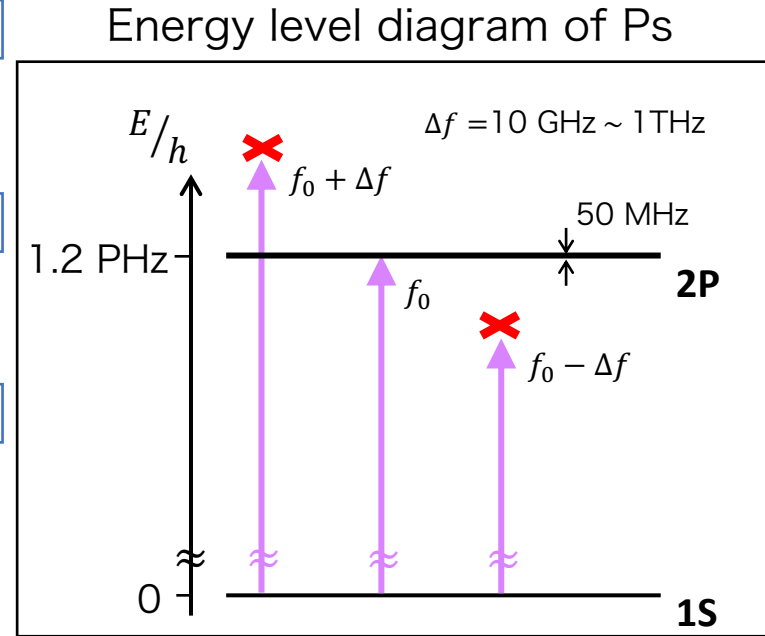
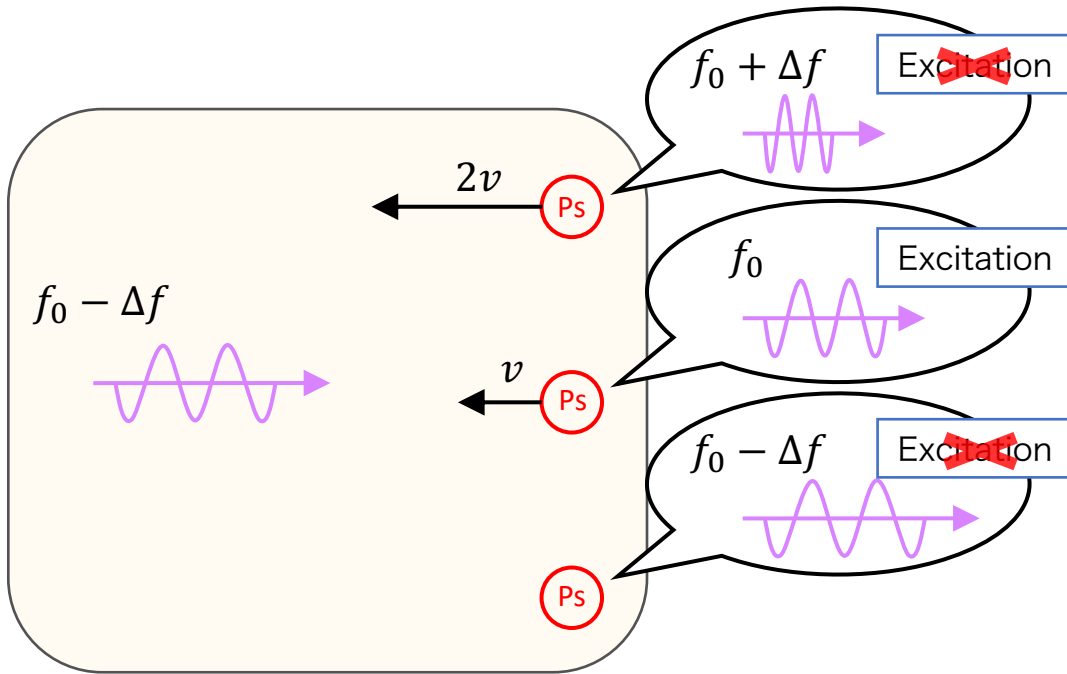
Spatial difference of 5 mm after flight of 800 ns

$\sigma$ -Ps decay time: 142 ns

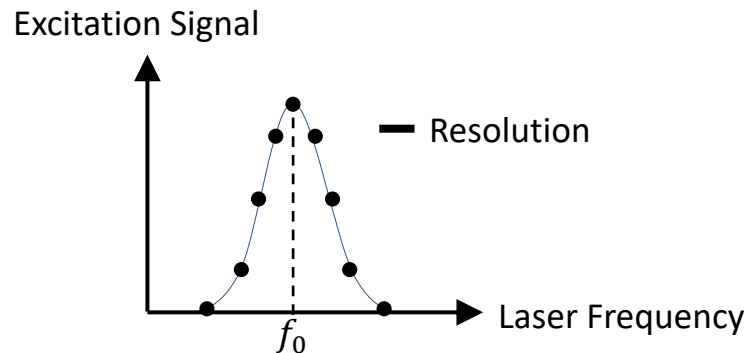
- Doppler broadening spectroscopy (DBS)
- Angular Correlation of Annihilation Radiation (ACAR)
- ★ Doppler Laser Spectroscopy

Doppler laser spectroscopy could be a suited method for cold Ps temperature measurement

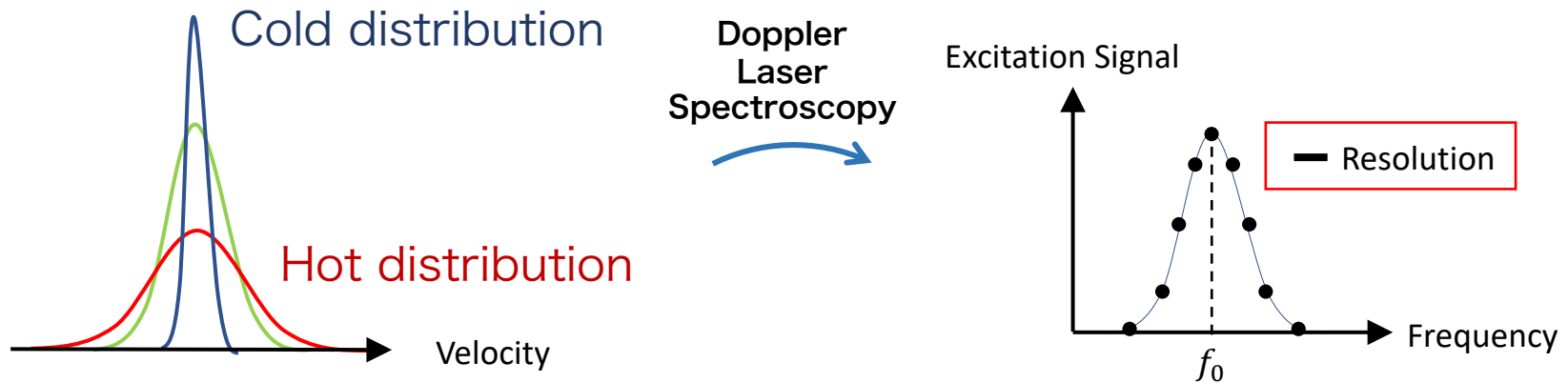
# Doppler Laser Spectroscopy of Ps



Doppler Laser Spectroscopy



# Resolution (given by Laser Linewidth) will be Tunable



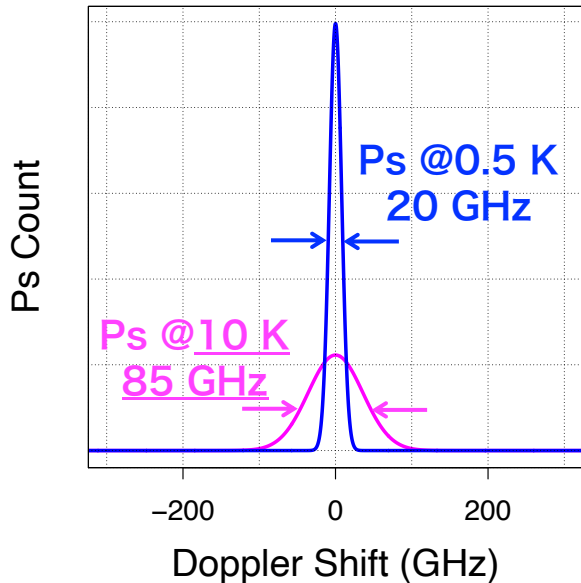
Resolution is determined by the linewidth of the probing laser (=how monochromatic the laser is)

With the laser technology of today, linewidth could be adjusted to a desired

Laser spectroscopy by a laser with an optimized linewidth could suit temperature measurement of cold Ps

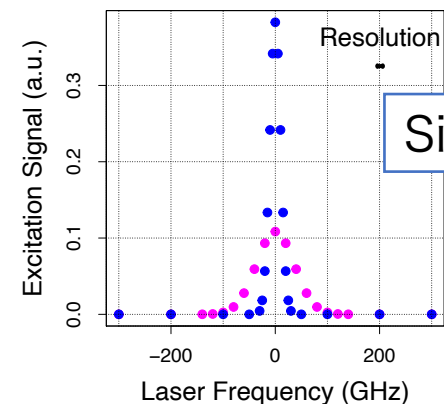
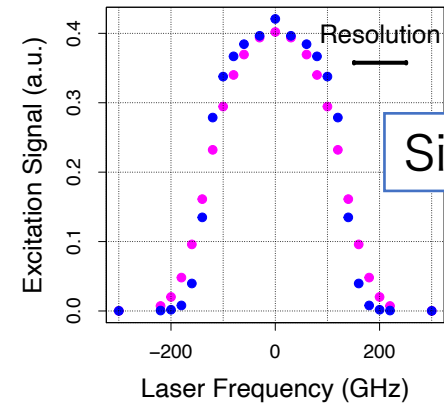
# How Narrow should the Linewidth be?

Ps doppler shift  
(Velocity distribution)



Typical available laser\*1  
Linewidth: **100 GHz**

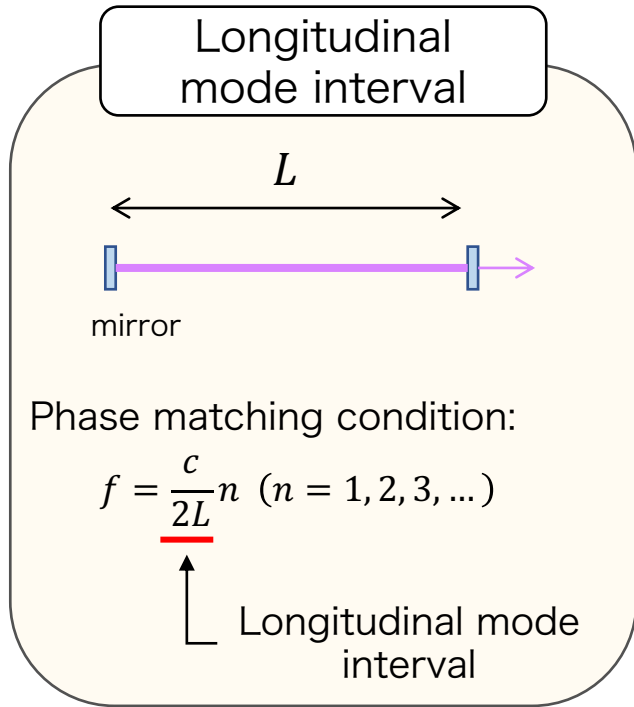
New laser  
Linewidth: **10 GHz**



Optimized linewidth for Ps @ 10 K  
would be 10 GHz

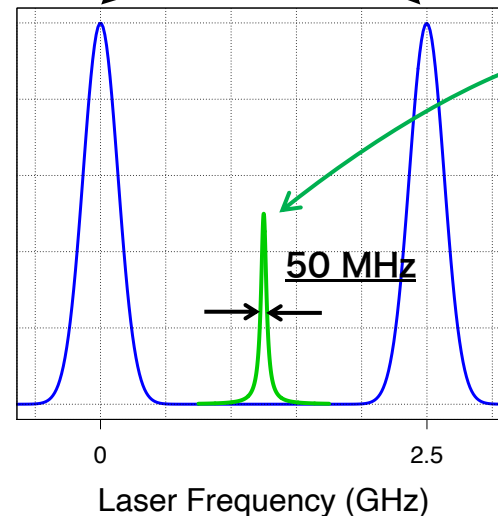
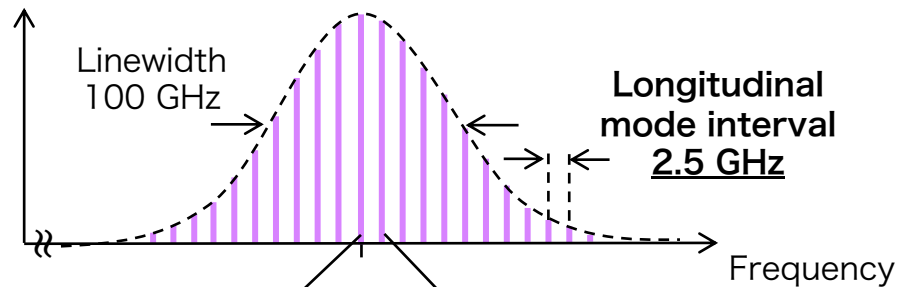
\*1 laser intensity: 500  $\mu\text{J} / (1 \text{ cm}^2 \times 2 \text{ ns})$  @243 nm

# Choosing an Appropriate Longitudinal Mode Interval



Longitudinal mode intervals should be narrow enough to be comparable to the natural width of Ps (50 MHz)

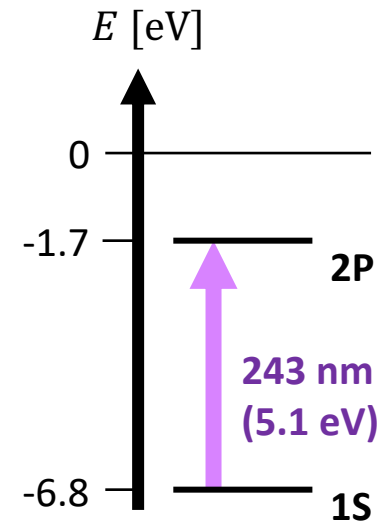
Frequency structure of a typical available laser



Spectrum is sparse for Ps

# Requirements for the Cold Ps Probing Laser

1. Optimized linewidth: around 10 GHz
2. Appropriate longitudinal mode interval: near 50 MHz
3. Wavelength: Ultraviolet (243 nm)
4. Adequate pulse intensity
5. Frequency stability

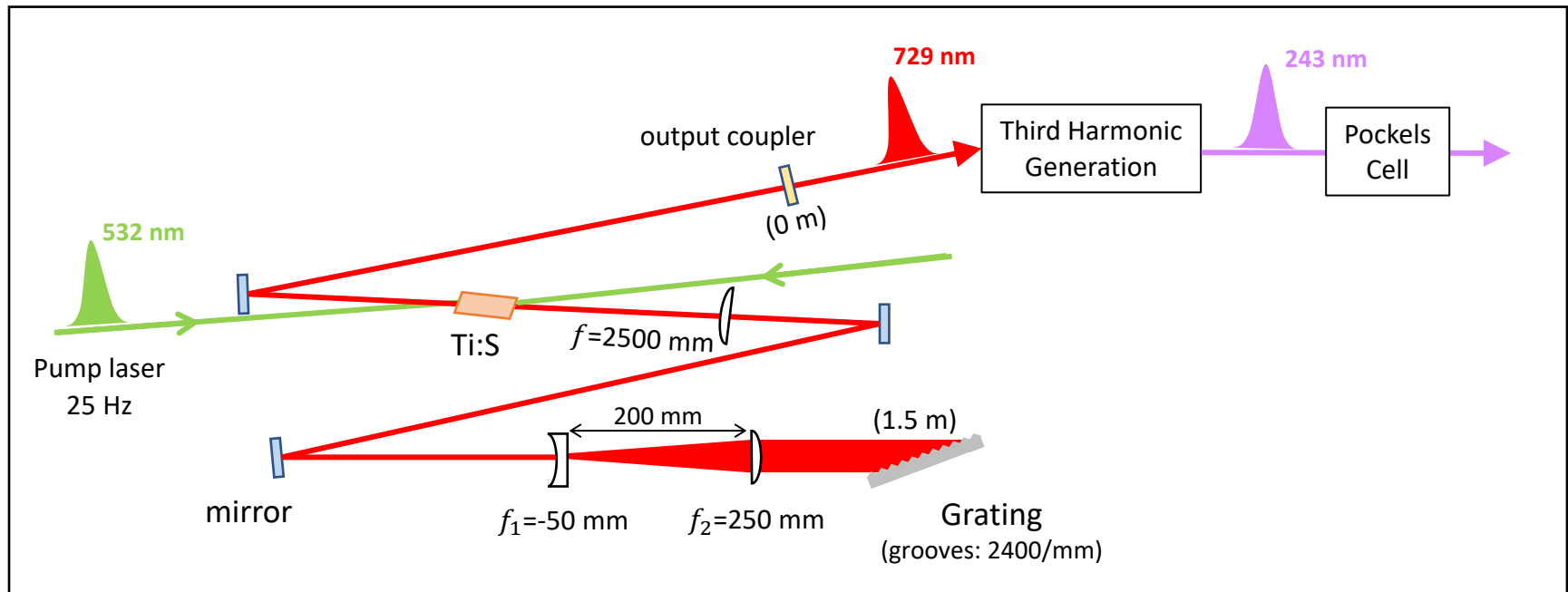


Energy level diagram of Ps

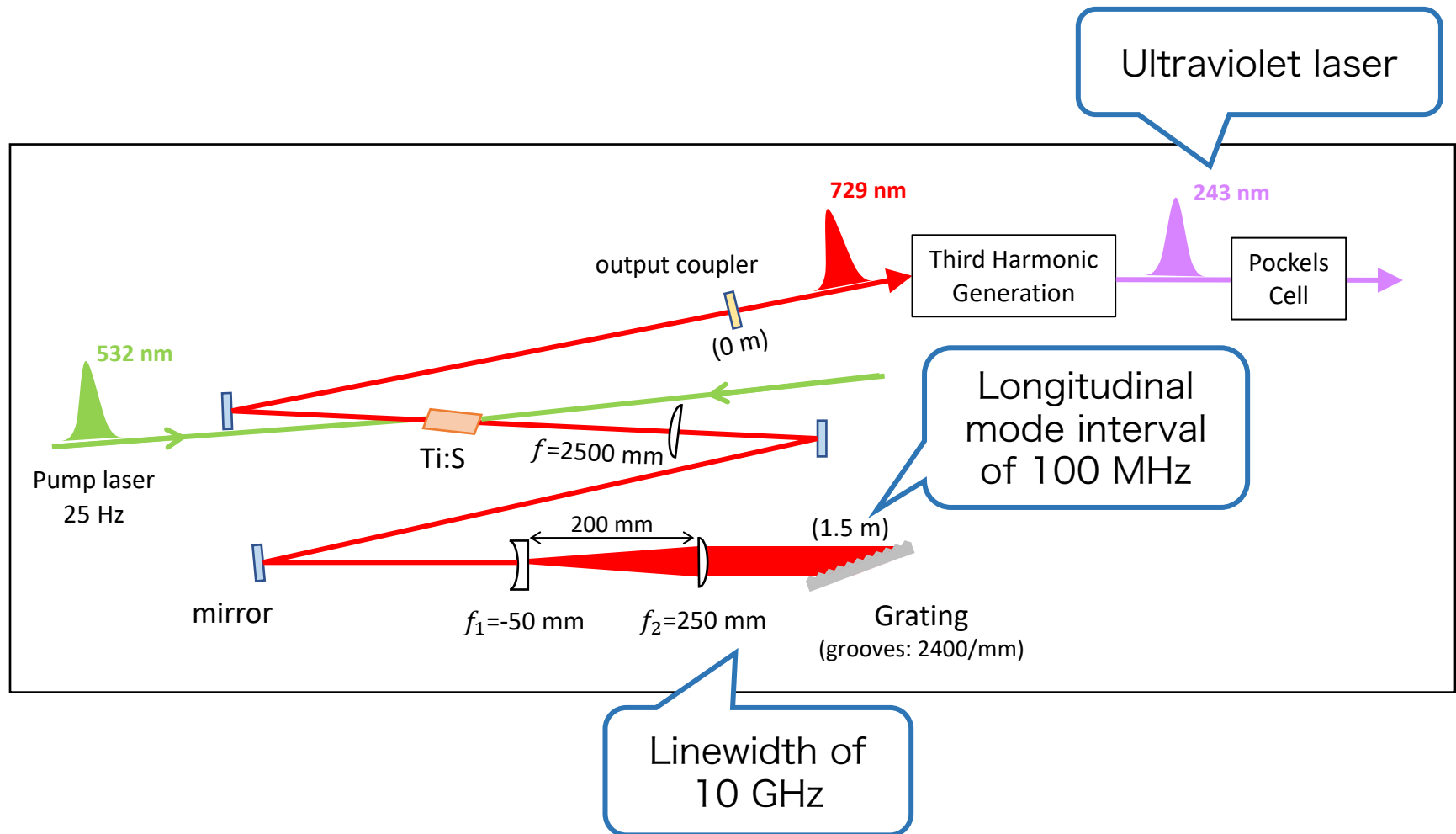
These requirements can not be achieved by commercially available lasers



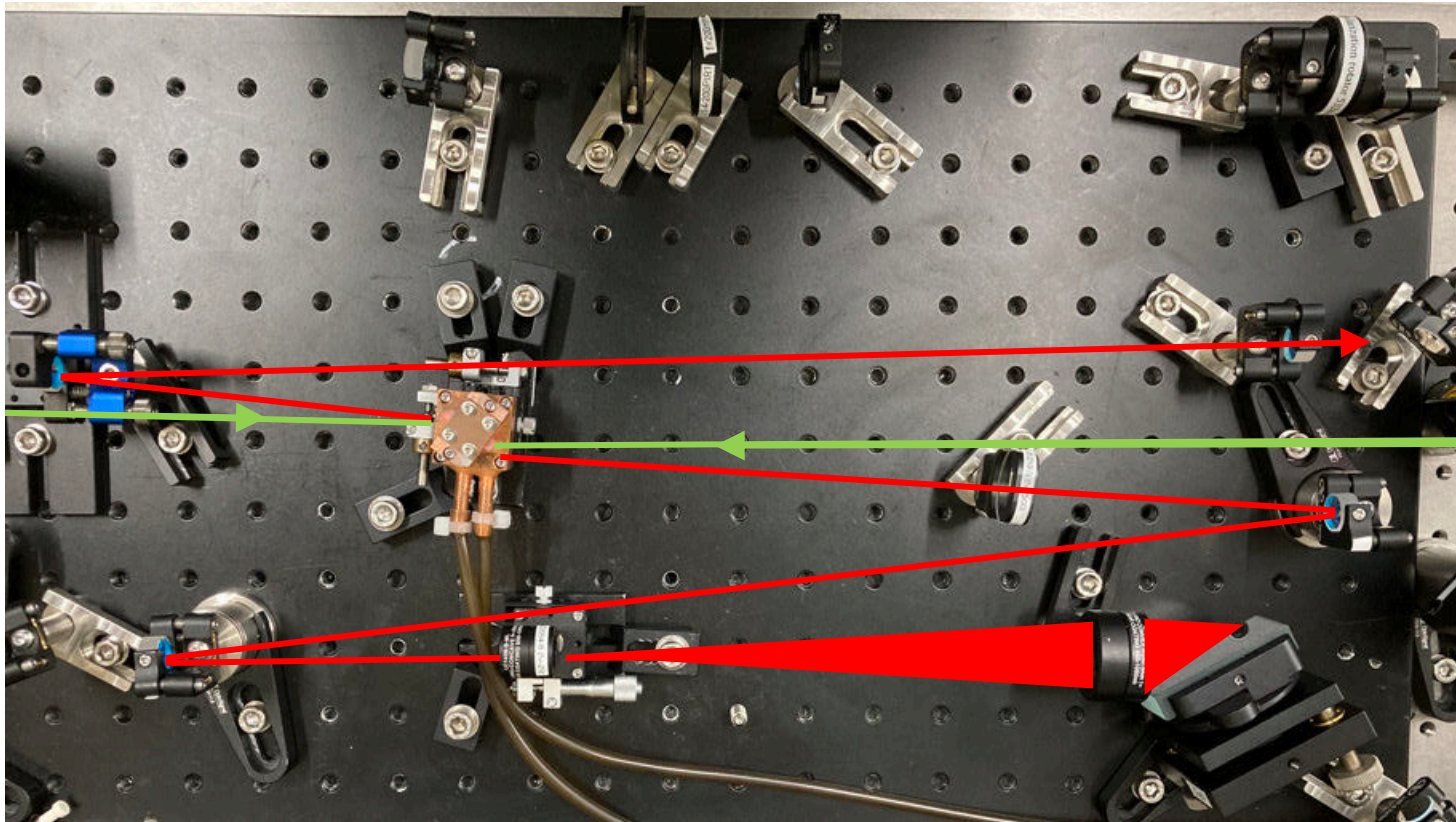
# The Newly Developed High-Resolution Probing Laser



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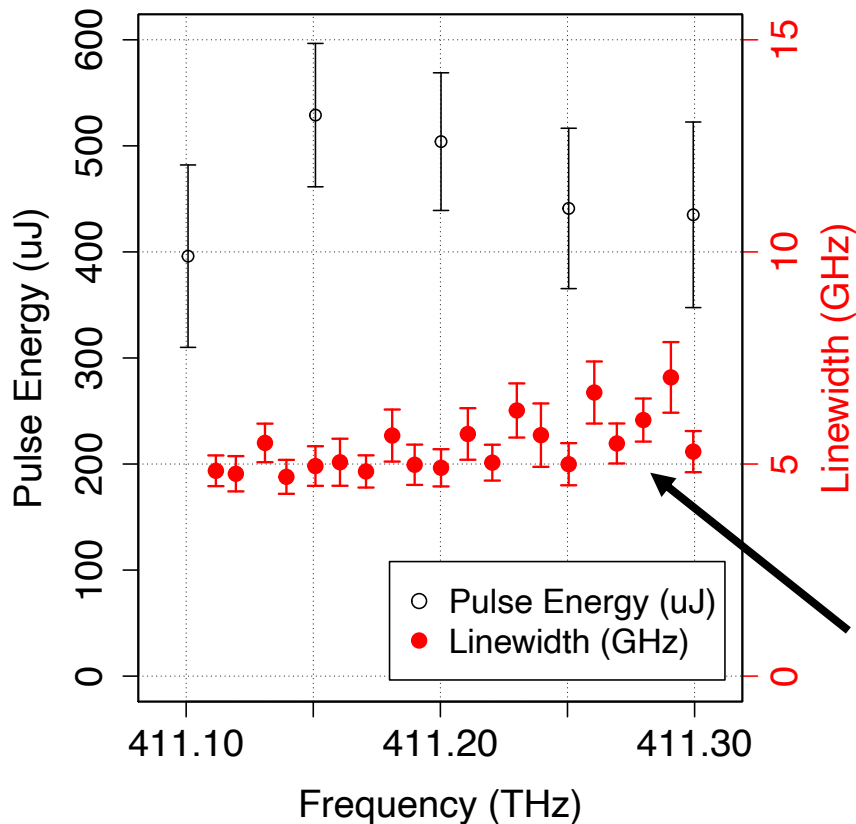


# The Newly Developed High-Resolution Probing Laser



# Linewidth

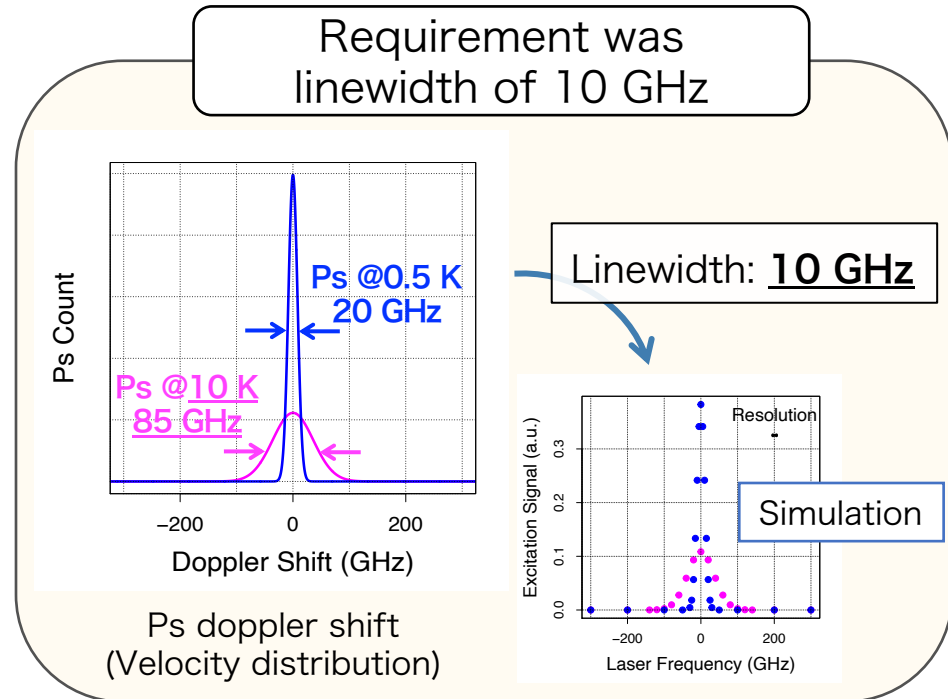
Linewidth measurement using Fizeau interferometer @729 nm



Linewidth of 6 GHz @729 nm

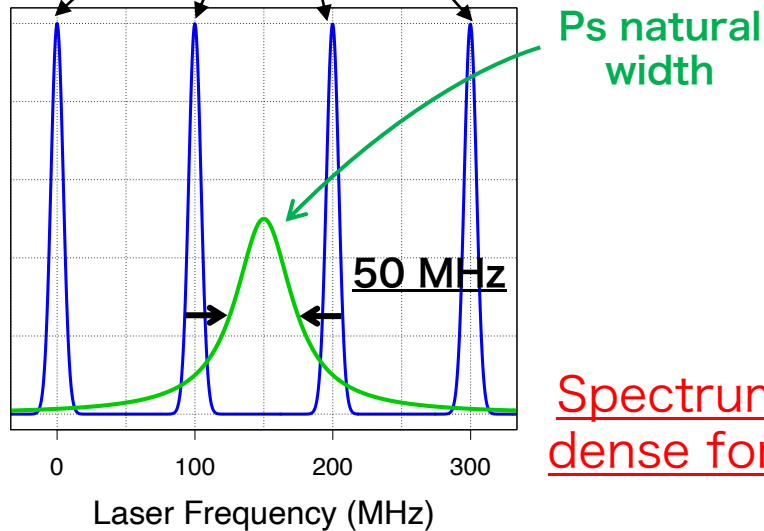
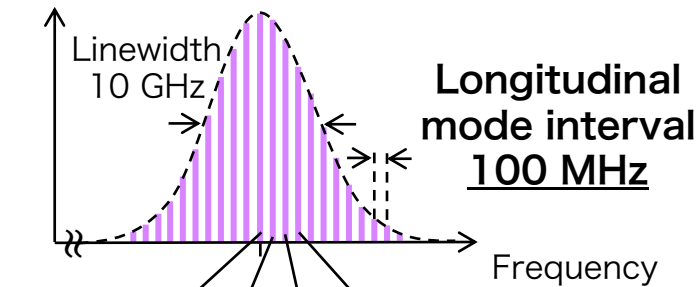
10 GHz @243 nm

Achieved the requirement of 10 GHz linewidth

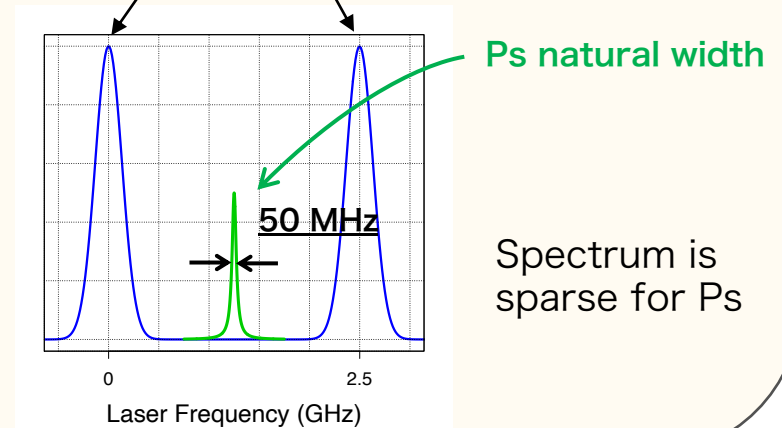
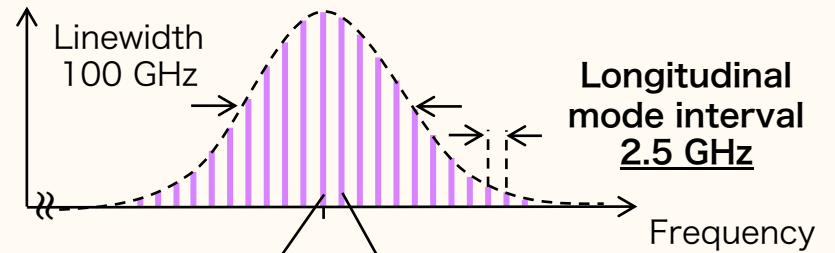


# Longitudinal Mode Interval

Frequency structure of the newly developed laser



Spectrum of typical available lasers would be sparse for Ps

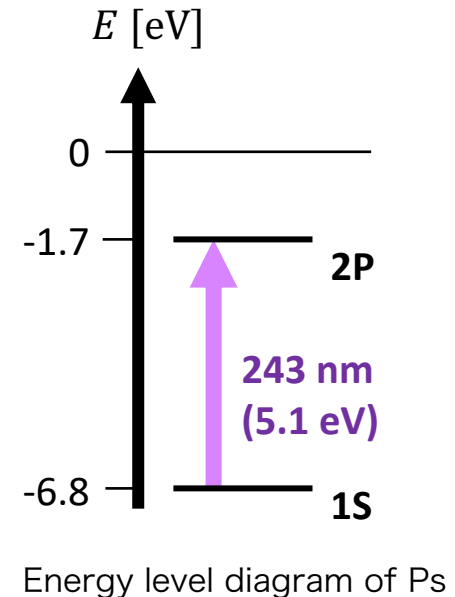


# Requirements for the Cold Ps Probing Laser

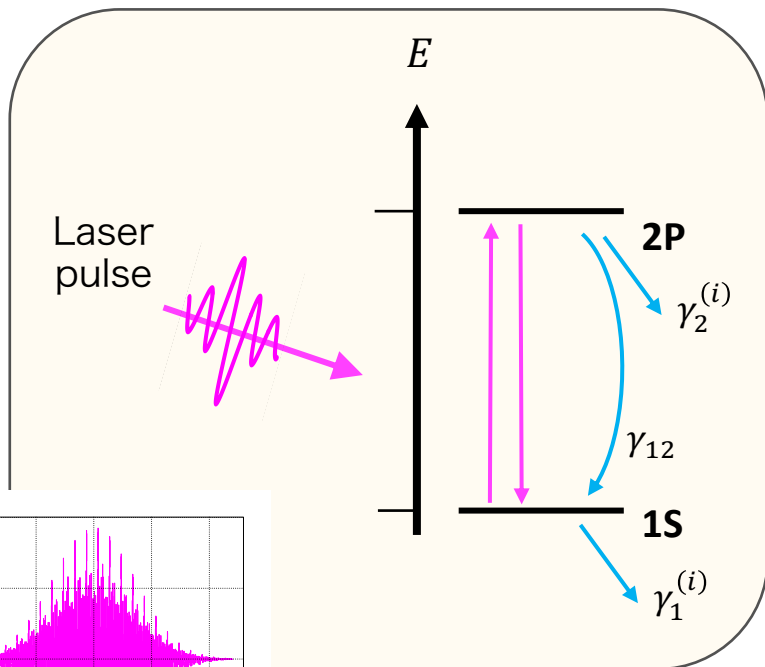
1. Optimized linewidth: around 10GHz
2. Appropriate longitudinal mode intervals: near 50 MHz
3. Wavelength: Ultraviolet (243 nm)
4. Adequate pulse intensity
5. Frequency stability

Small fluctuation to endure long measurements

Excite enough Ps to 2P state



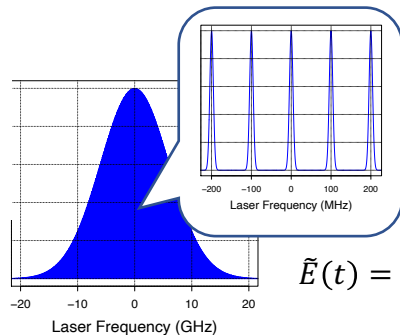
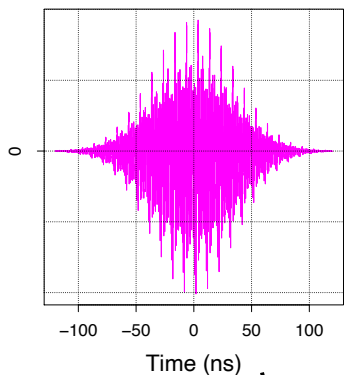
# Numerical Simulation of Laser Induced Excitation



Track the density matrix elements

$$\begin{cases} \frac{d\rho_{11}}{dt} = -\frac{i}{2}(\Omega\tilde{\rho}_{12} - \Omega^*\tilde{\rho}_{21}) - \sum_i \gamma_1^{(i)}\rho_{11} + \gamma_{12}\rho_{22} \\ \frac{d\tilde{\rho}_{12}}{dt} = -\left(\frac{1}{T_2} - i\Delta\omega\right)\tilde{\rho}_{12} - \frac{i}{2}\Omega^*(\rho_{11} - \rho_{22}) \\ \frac{d\tilde{\rho}_{21}}{dt} = -\left(\frac{1}{T_2} + i\Delta\omega\right)\tilde{\rho}_{21} + \frac{i}{2}\Omega(\rho_{11} - \rho_{22}) \\ \frac{d\rho_{22}}{dt} = \frac{i}{2}(\Omega\tilde{\rho}_{12} - \Omega^*\tilde{\rho}_{21}) - \sum_i \gamma_2^{(i)}\rho_{22} - \gamma_{12}\rho_{22} \end{cases}$$

Electric Field

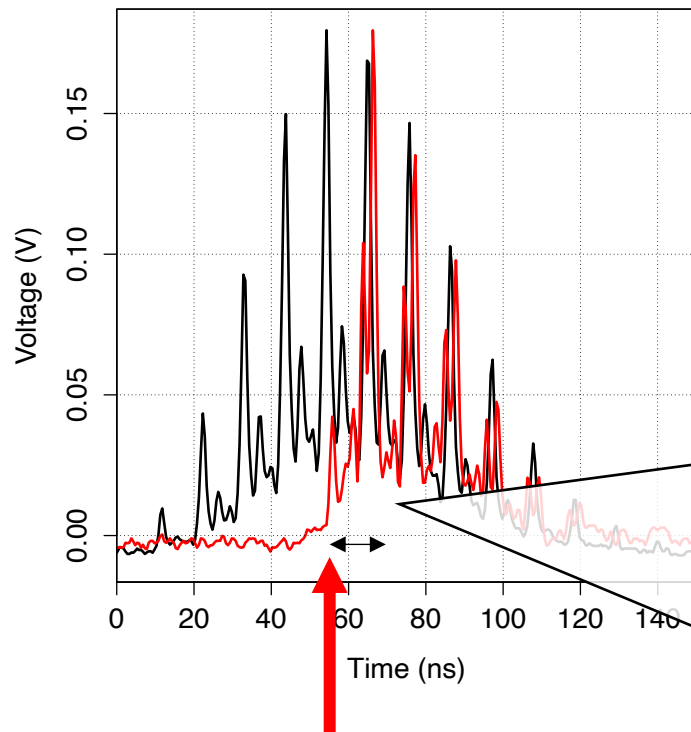


$\mathcal{F}$

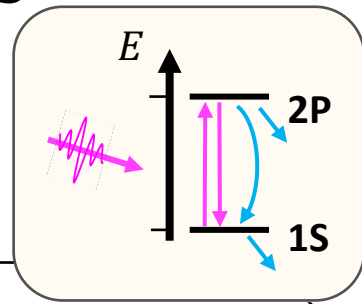
$$\tilde{E}(t) = \sqrt{2\pi}A_0\sigma_m \exp\left[-\frac{\sigma_m^2 t^2}{2}\right] e^{-i(\omega_c + \delta\omega)t} \sum_k \exp\left[-\frac{(k\omega_{\text{rep}} + \delta\omega)^2}{2\sigma^2}\right] e^{-i(k\omega_{\text{rep}}t - \phi(k,t))}$$

# Adequate Pulse Intensity to Excite Ps to 2P State

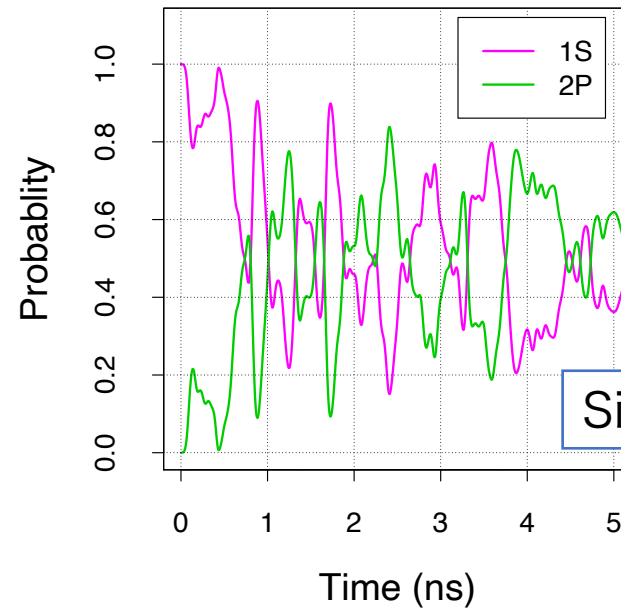
Pulse waveform from a photodetector



Cut pulse using Pockels cell



Time evolution of density matrix

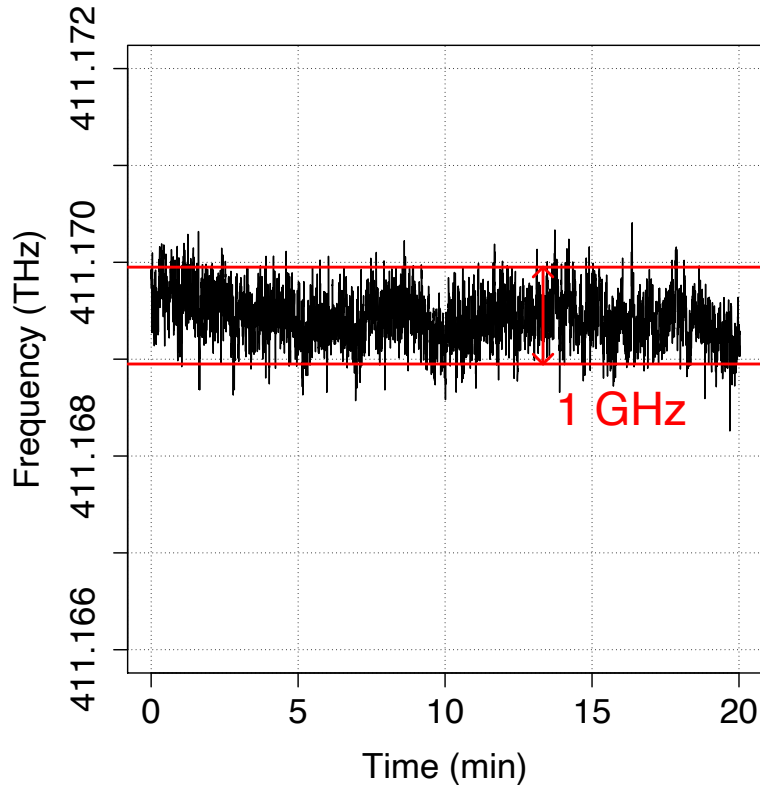


\*1 laser intensity:  $25 \text{ uJ} / (1 \text{ cm}^2 \times 60 \text{ ns})$



# Frequency Stability

Frequency measurement  
using a wavemeter @729 nm



20 min frequency stability of  
1 GHz @729 nm  
(3 GHz @243 nm)

Sufficient frequency stability  
to endure long measurements

# Conclusion and Future Prospects

## Conclusion

- Developed a new high-resolution probing laser.
- Achieved an optimized linewidth of 10 GHz which could probe Ps at 10 K or below.
- Narrow longitudinal mode interval of 100 MHz was achieved to efficiently excite Ps.
- The new laser has adequate pulse intensity to pump enough Ps to excitation state.

## Future prospects

- Preparing for spectroscopy measurements using the new laser.
- Application of utilizing the new laser to observe proof of laser cooling.