

Evaluation of energy-resolved spin polarization of surface electrons by spin-polarized positronium time-of-flight method (SP-PsTOF)

M. Maekawa, S. Sakai, A. Miyashita and A. Kawasuso

***Quantum Beam Science Research Directorate,
Takasaki institute,
National Institutes for Quantum and Radiological
Science and Technology***

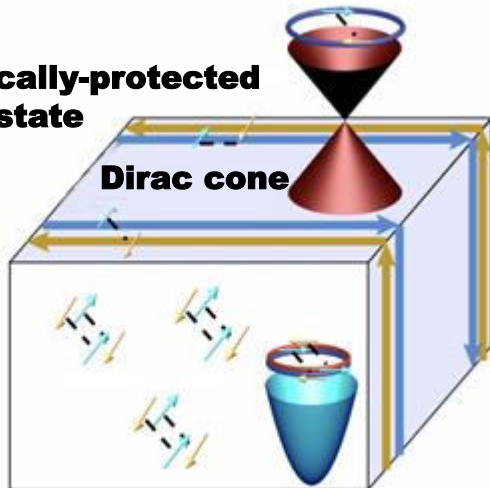
Topmost spin phenomena

- ✓ **Novel quantum spin phenomena on the topmost surface are important for realizing next-generation devices.**
- ✓ **To understand these phenomena, the spin polarization of the topmost layer must be evaluated correctly.**

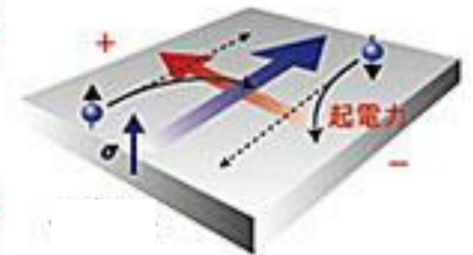
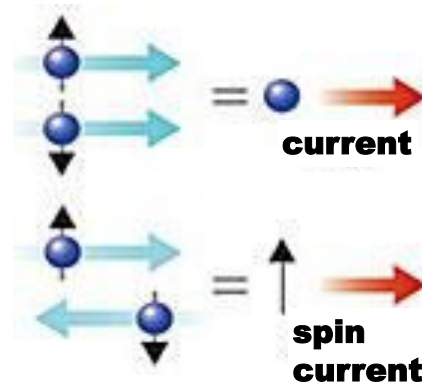
Spin phenomena on the topmost surface

ex. Topological insulator

topologically-protected
surface state



ex. spin-Hall effect



Surface spin evaluation

- ✓ There are several methods to evaluate surface spins.
- ✓ It is very difficult to evaluate only such the spins.

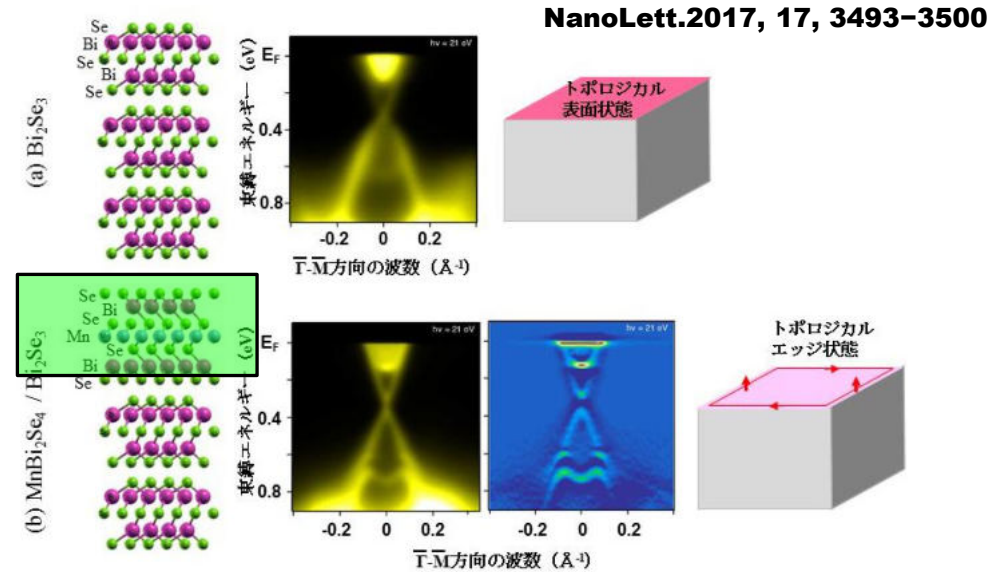
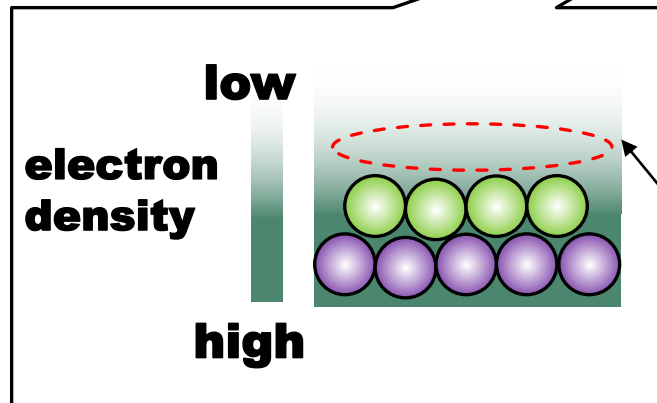
Commonly used techniques:

S-ARPES

Spin- and Angle-Resolved
Photoemission Spectroscopy

DP-XMCD

Depth-profiling X-ray Magnetic
Circular Dichroism



These results include information not only from the surface electrons but also from the several bulk layer.

Novel quantum spin phenomena occurs here

Positronium spectroscopy can provide such information

Advantages of positronium spectroscopy:

1. Ps is formed only at the **vacuum side of the surface.**
2. Ps annihilation changes depending on the **direction of electron spin.** (p-Ps and o-Ps)
3. Ps emission affected by the **state of electrons.**

1. Ps formation only at surface

① incident of positrons

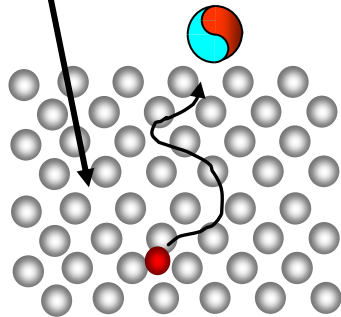


④ annihilation



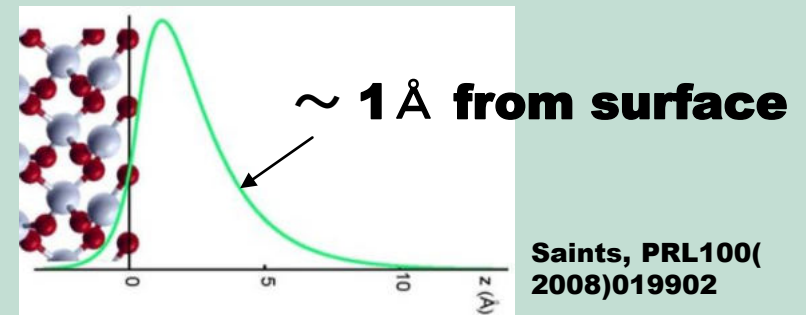
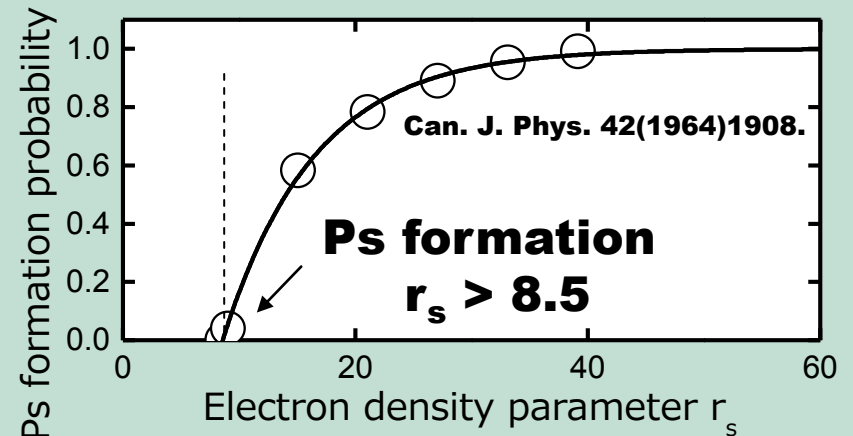
⑤ Ps emission

③ Ps formation with surface electrons



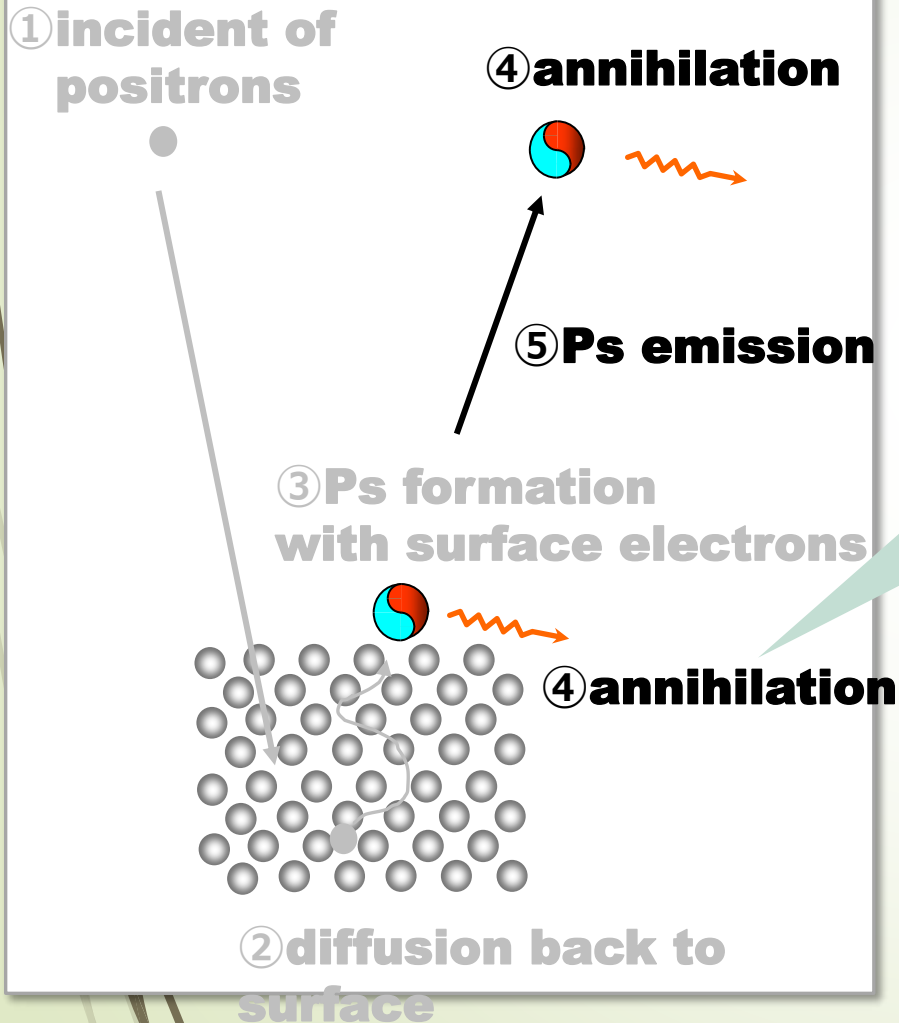
② diffusion back to surface

Ps can form only at surface



Ps observation = surface electron observation

2. Ps state changes by electron spin



Ps state changed by electron spin

spin parallel



ortho-Ps

spin anti-parallel



para-Ps

3 photon decay
($\sim 511\text{keV} \times 3$)

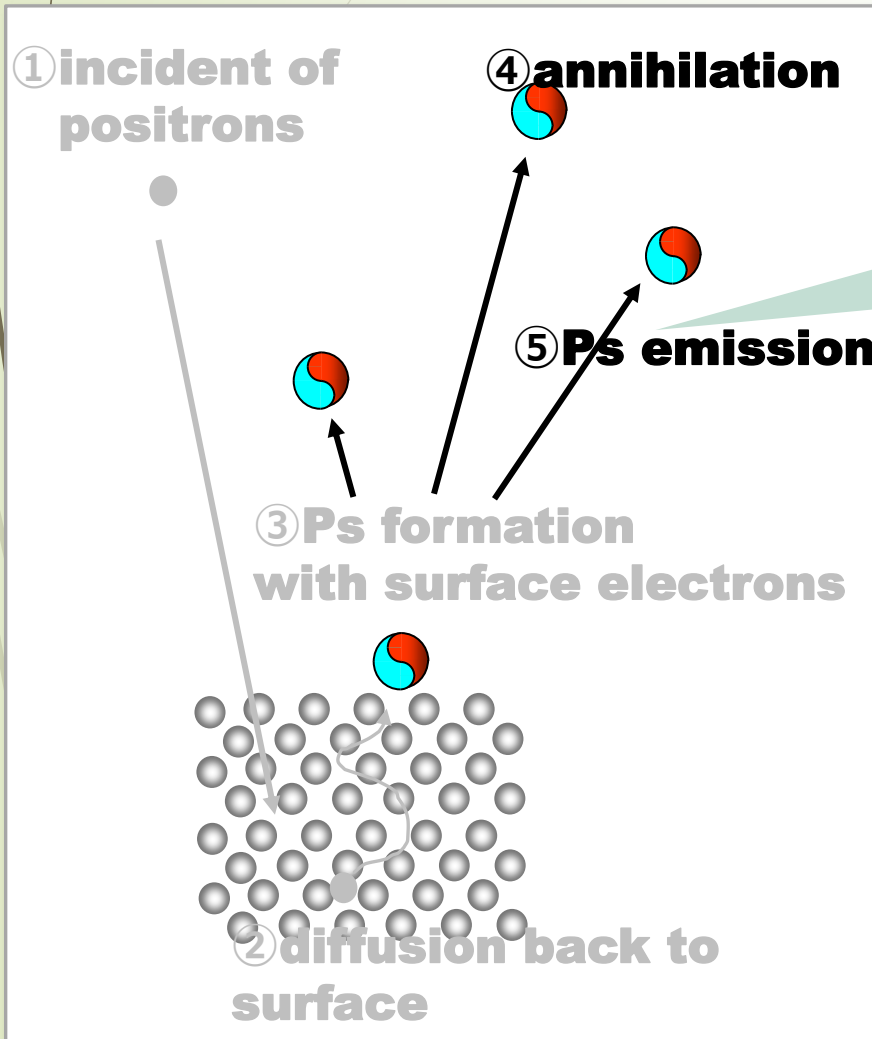
2 photon decay
($511\text{keV} \times 2$)

lifetime **142ns**

lifetime **125ps**

**Determination of electron spins
(i.e. spin polarization)**

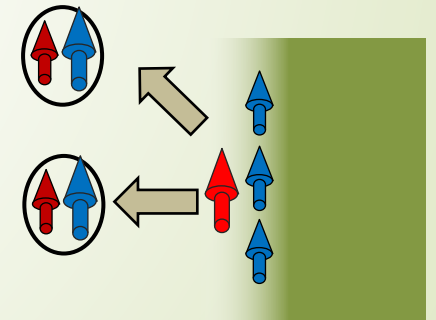
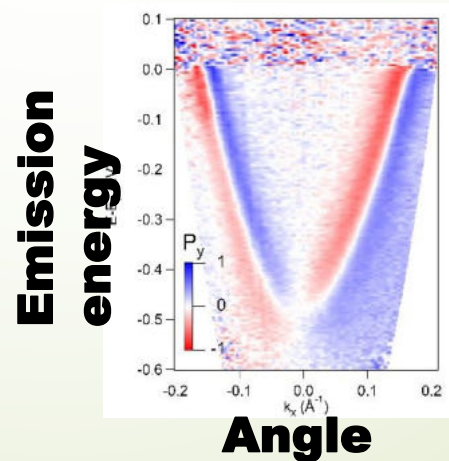
3. Ps emission affected by electron state



Ps emission energy and angles are affected by electron state

Ps emission energy spectrum

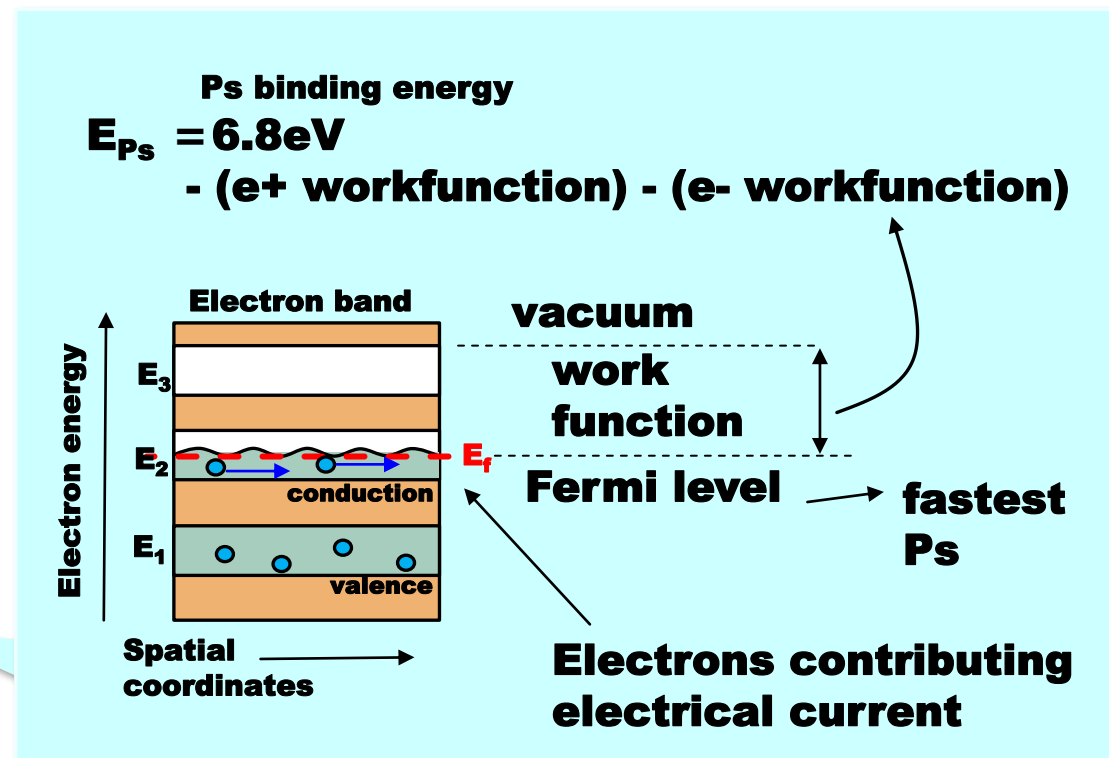
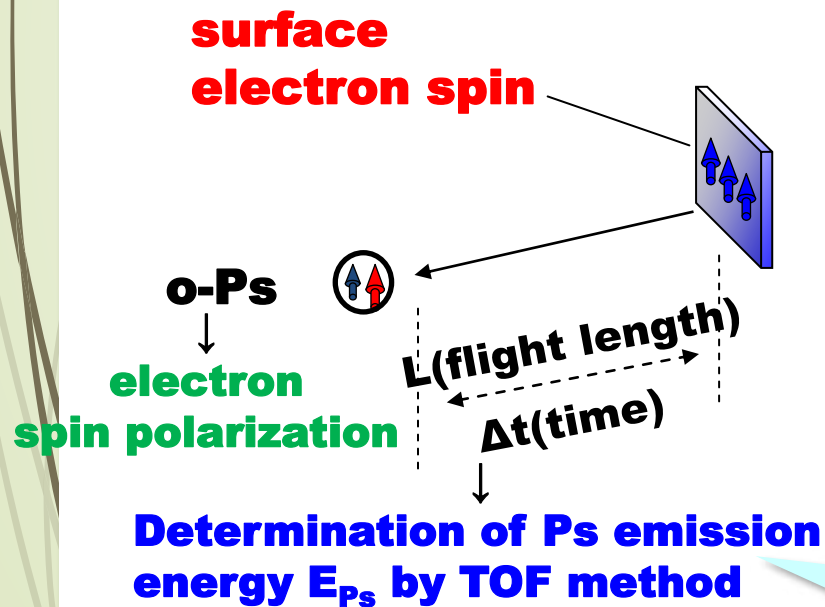
$$\frac{d(N(\theta\varphi))}{d\varepsilon_{\perp}} \propto \frac{\alpha}{\sqrt{\varepsilon_{\perp}}} \int \gamma \left[\tan(\theta) - \frac{|k_{\parallel}|}{q} \right] |\nabla_{\perp}(\varepsilon\mathbf{k})|^{-1} dk_{\parallel}$$



Spin-polarized electron band structure might be obtained !

Energy-resolved Ps spectroscopy

- ✓ Energy of emitted Ps can be determined by **Ps time-of-flight(TOF)** .
- ✓ The spin-polarized electron density of states associated only with the topmost layer of metals will be obtained.



PsTOF apparatus

Timing signal

beam pulsing (high efficiency),
spin depolarization

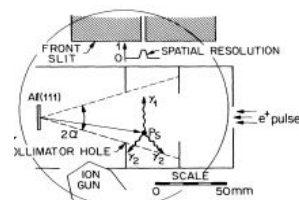
secondary electron
spin pol.

Positron source

**Large-scale
facility**

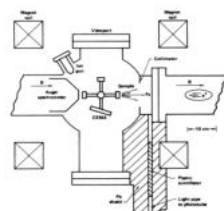
**high
intensity**

Bell lab.



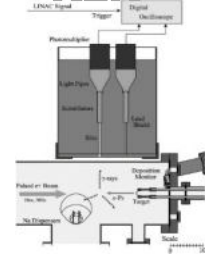
PRB34(1986)3069

LLNL



PRL51(1983)1085

KEK



Surf. Sci. 641(2015)68.

**Toront univ.
→ NEPOMUC**

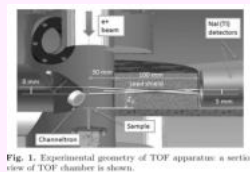
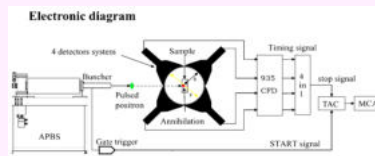


Fig. 1. Experimental geometry of TOF apparatus: a section view of TOF chamber is shown.

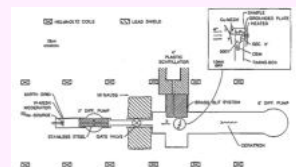
RI source

**low intensity,
spin pol.**

IHEP



London univ, ETHZ



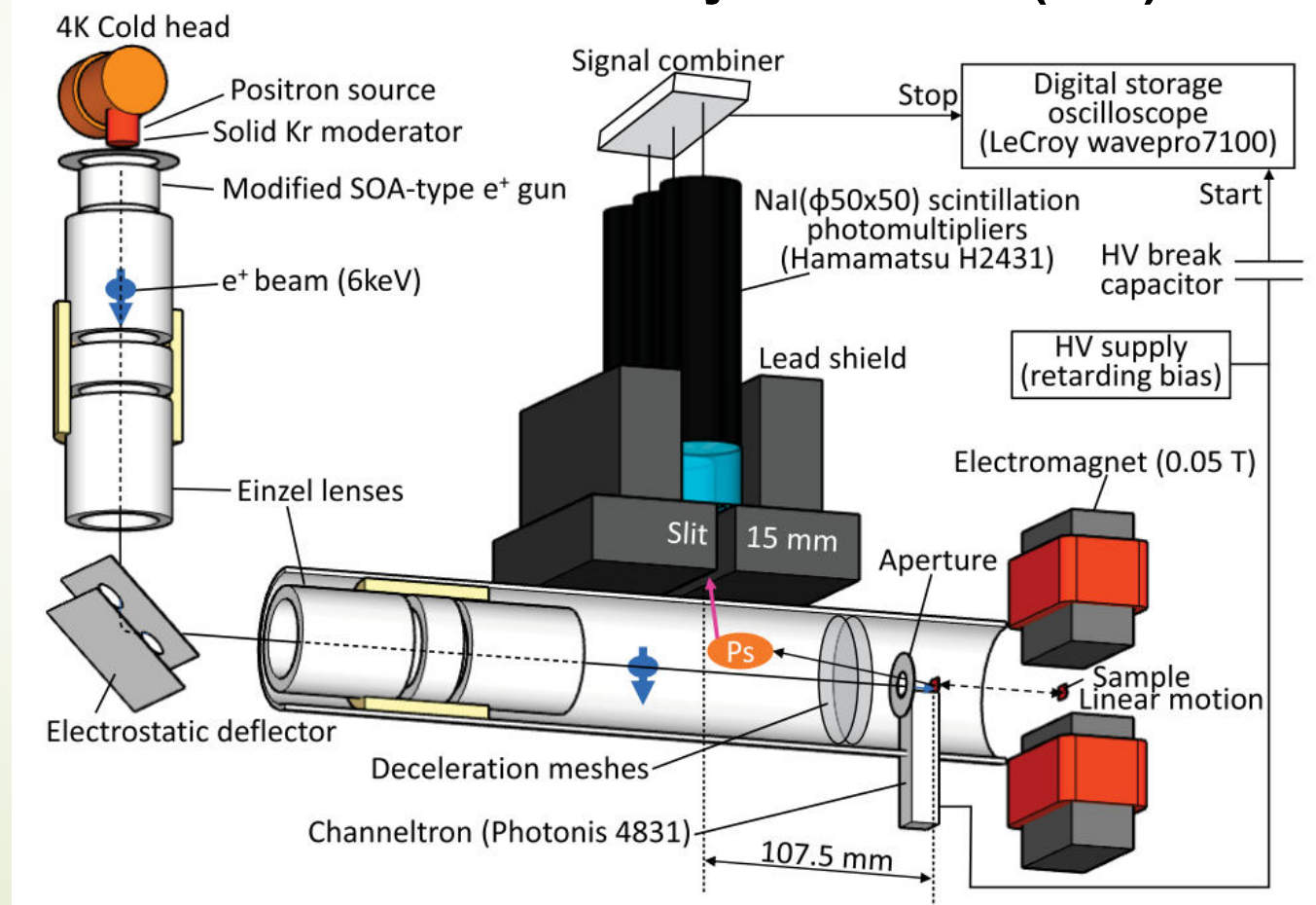
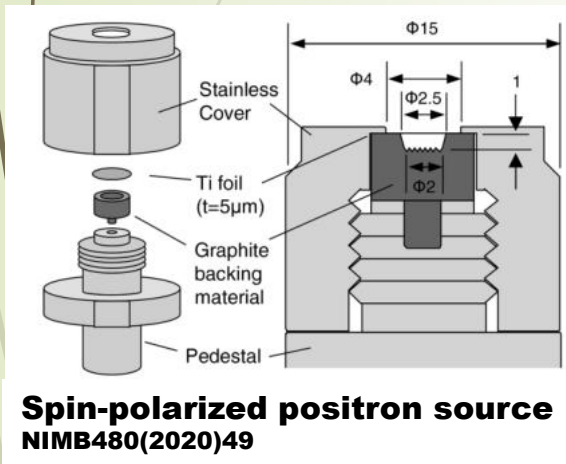
no one

Harder than we thought...(extremely low counting rate)

Spin-polarized Ps-TOF (SP-PsTOF)

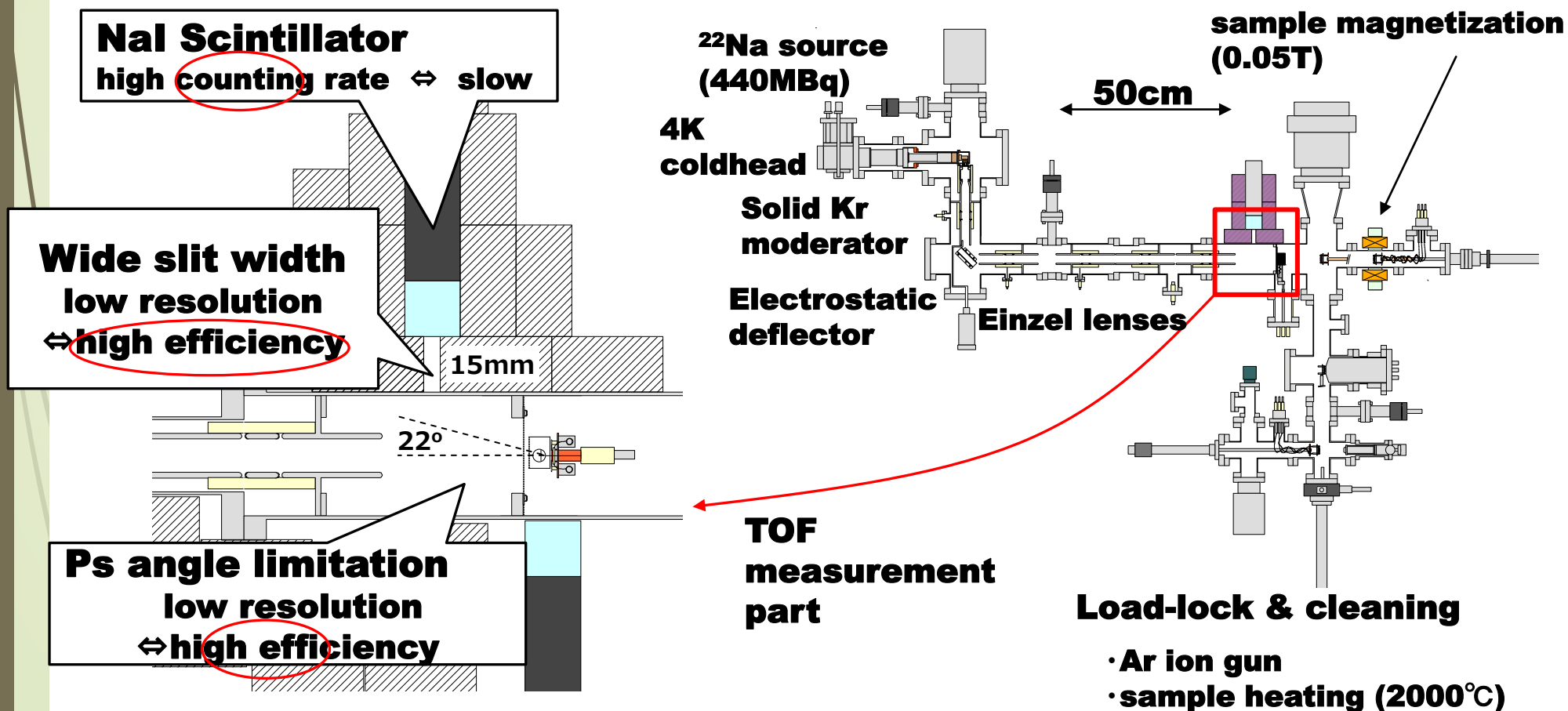
✓ Construction a Ps-TOF apparatus using spin-polarized positron beam.

Phys. Rev. Lett.126(2021)186401



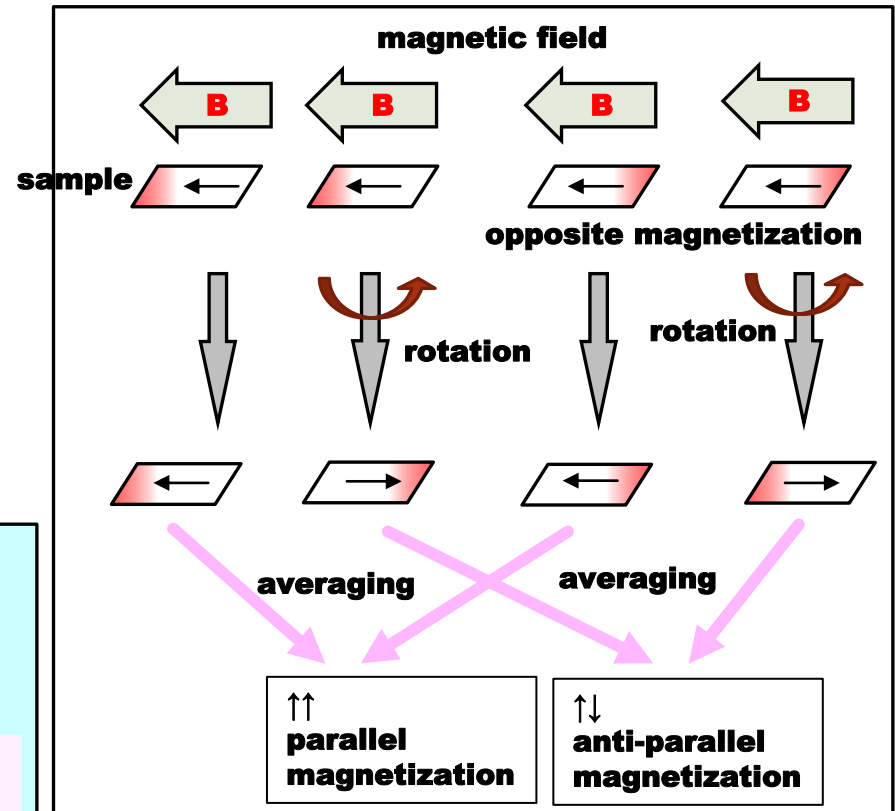
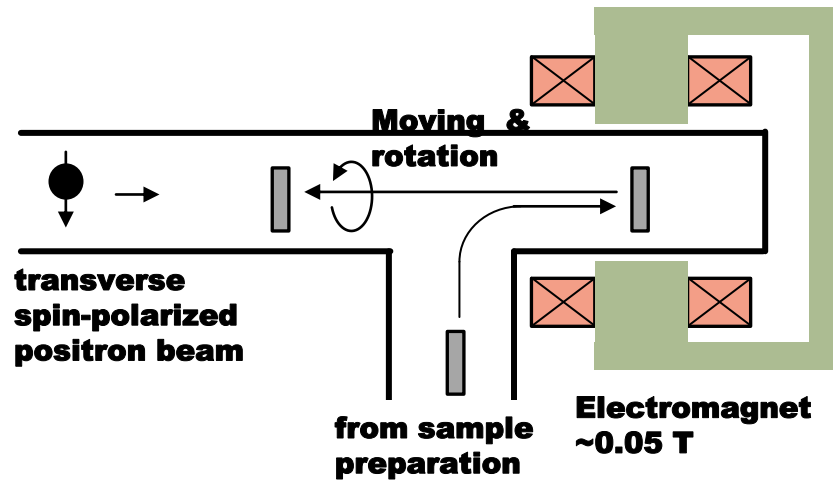
TOF detector

✓ Keeping counting rate is the highest priority.



Sample magnetization

- ✓ Sample magnetization by external magnetic field of 0.05T.
- ✓ The spectrum asymmetry was obtained from the spin parallel / antiparallel TOF spectra.



Cancellation of position dependency

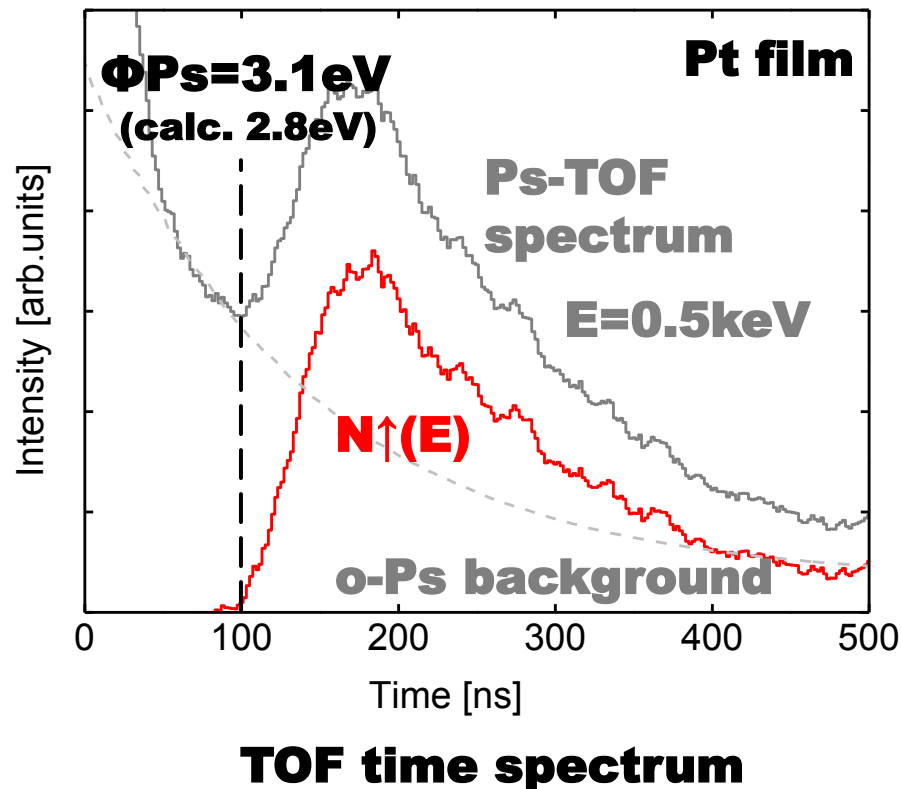
TOF spectrum { Neg. field : $I^\uparrow(E)$
Pos. field : $I^\downarrow(E)$

Asymmetry

$$A^{oPs} = \frac{I^\uparrow(E) - I^\downarrow(E)}{I^\uparrow(E) + I^\downarrow(E)} \propto P_-(E)$$

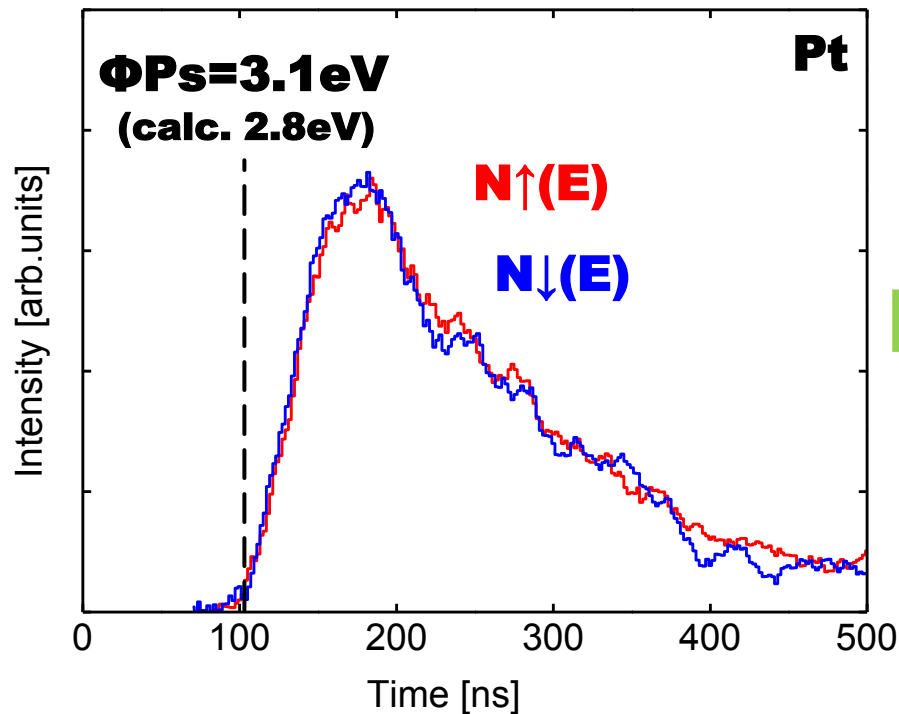
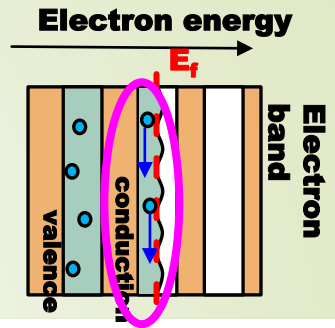
SP-PsTOF result (Pt film: non-magnetic)

- ✓ The differential Ps-TOF spectrum was measured for the non-magnetic Pt thin film.

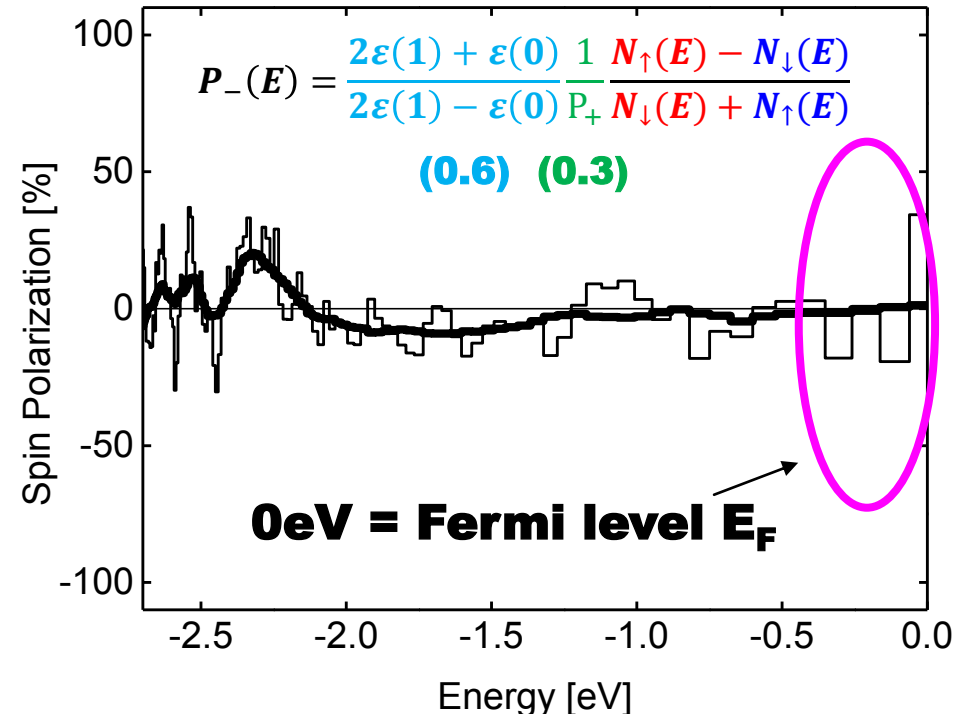


SP-PsTOF result (Pt film: non-magnetic)

- ✓ Difference PsTOF by alternating magnetization.
TOF → Ps energy (adjusted to $E_f=0$)
Intensity asymmetry → electron spin polarization
- ✓ Spin polarization $\sim 0\%$ (non-magnetic)



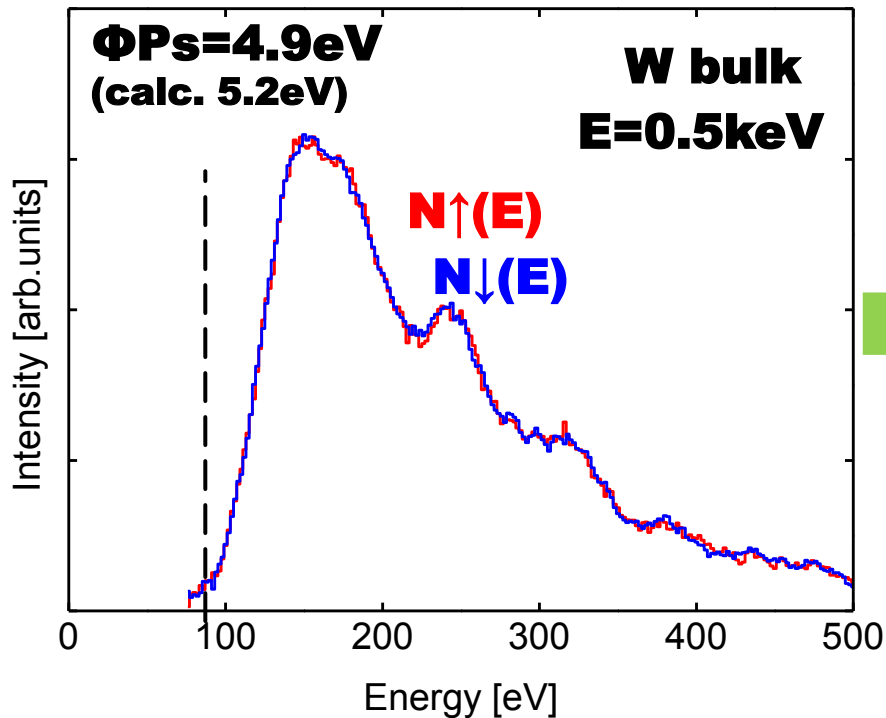
PsTOF time spectrum



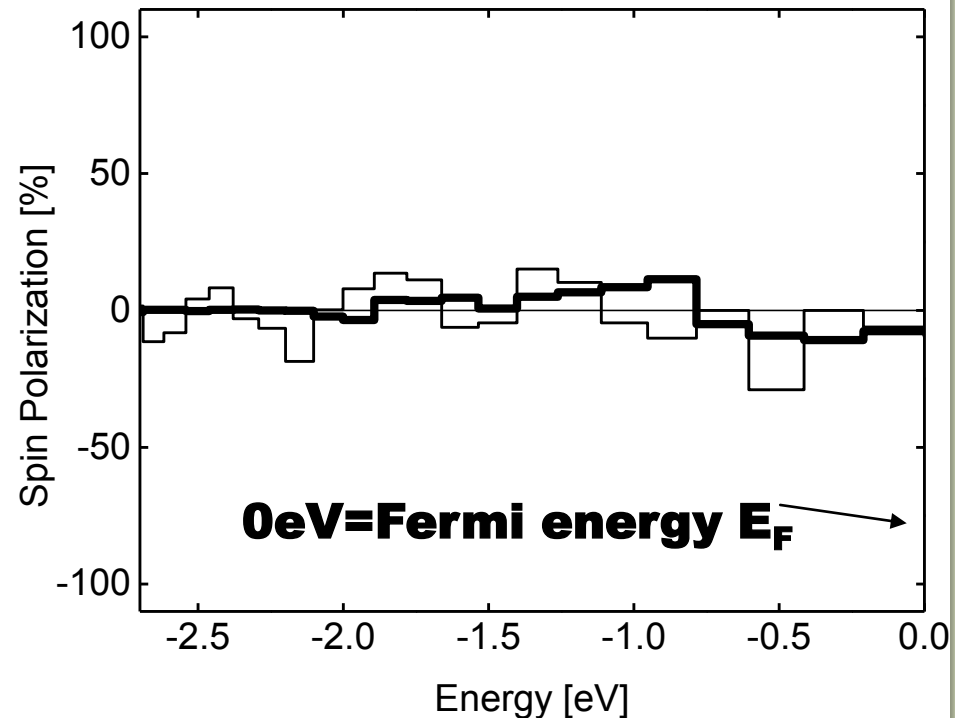
Energy-resolved topmost electron spin polarization

SP-PsTOF result (W bulk: non-magnetic)

- ✓ The differential Ps-TOF spectrum was measured for the non-magnetic W bulk crystal.
- ✓ Spin polarization $\sim 0\%$ (non-magnetic)



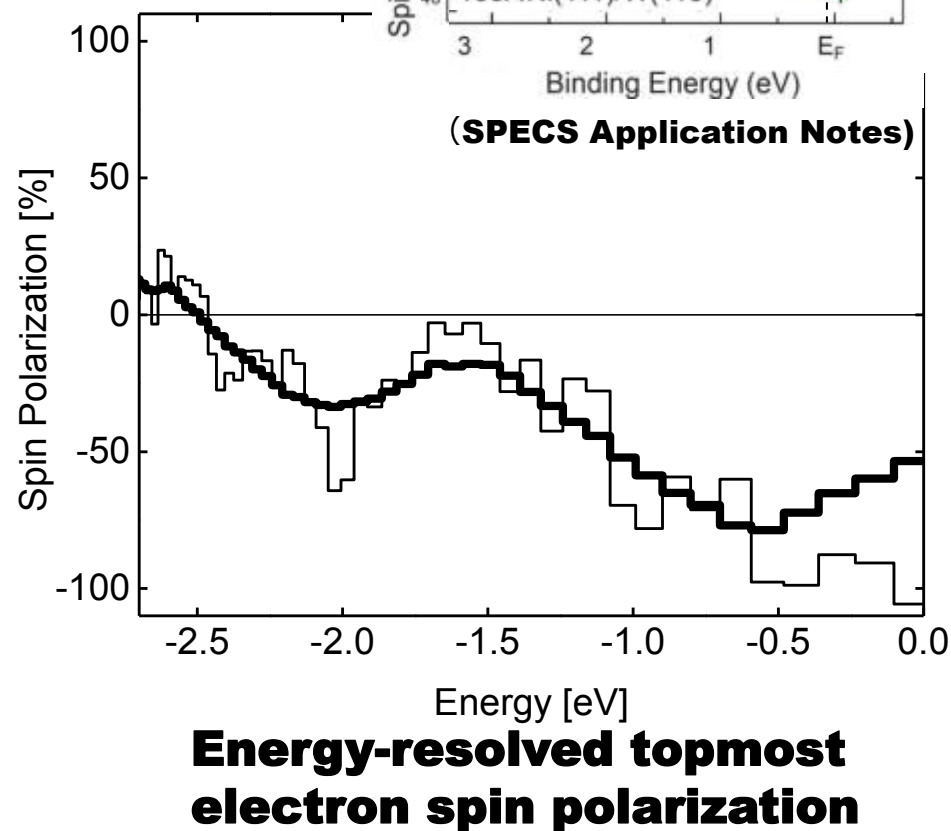
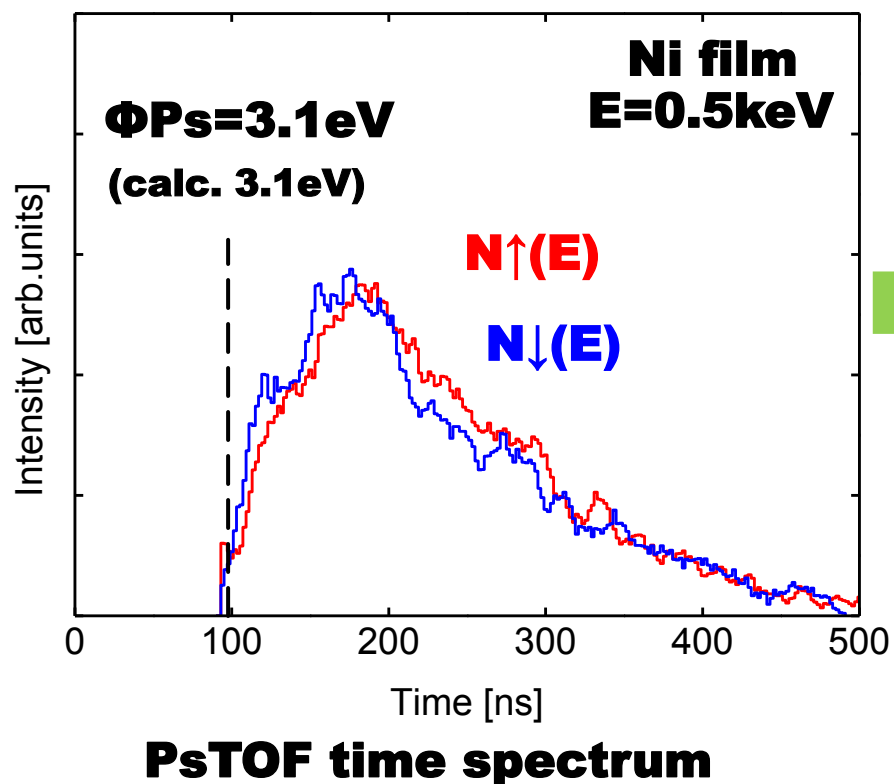
PsTOF time spectrum



**Energy-resolved topmost
electron spin polarization**

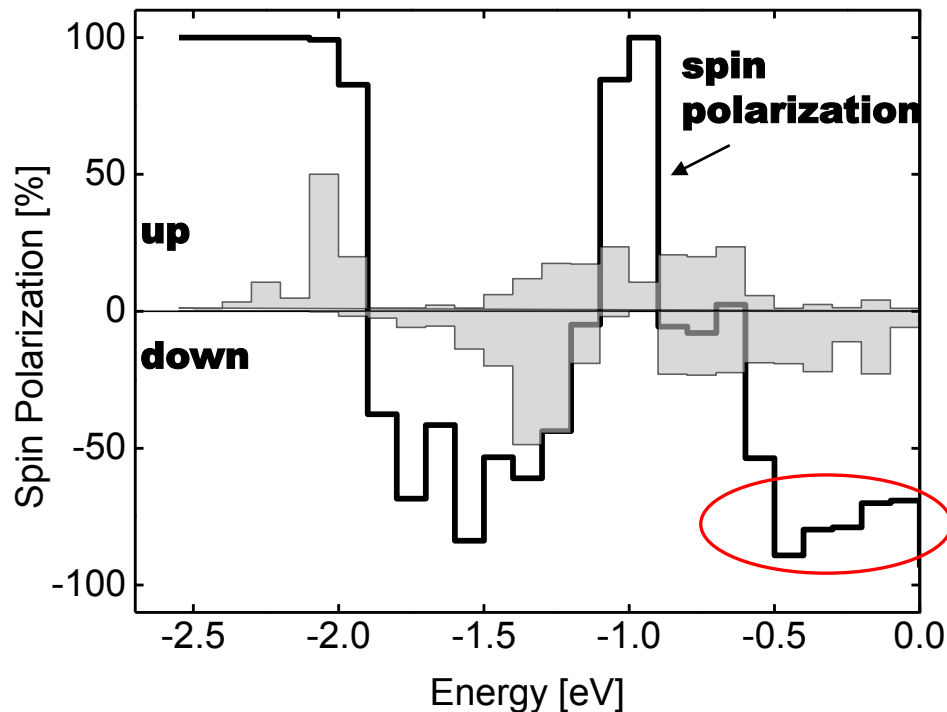
SP-PsTOF result (Ni film: magnetic)

- ✓ Clear negative polarization is appeared near the Fermi level.
- ✓ Same tendency as other methods

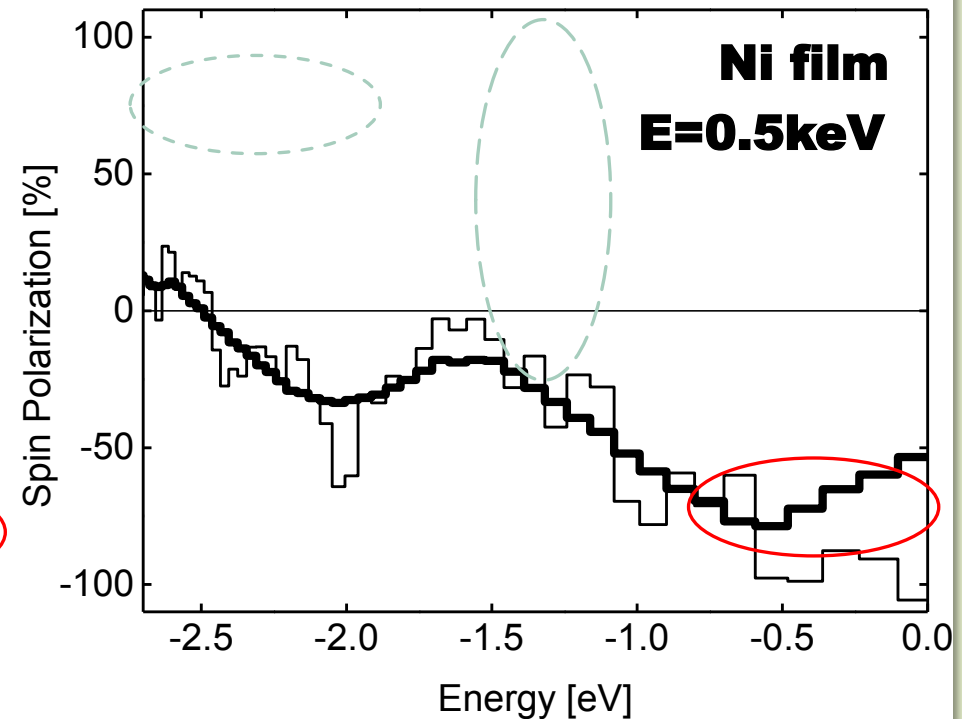


SP-PsTOF result (Ni film: magnetic)

- ✓ Comparison with first-principles calculation.
- ✓ Experimental and calculation show basically the same tendency, but they seem not to be exactly the same.



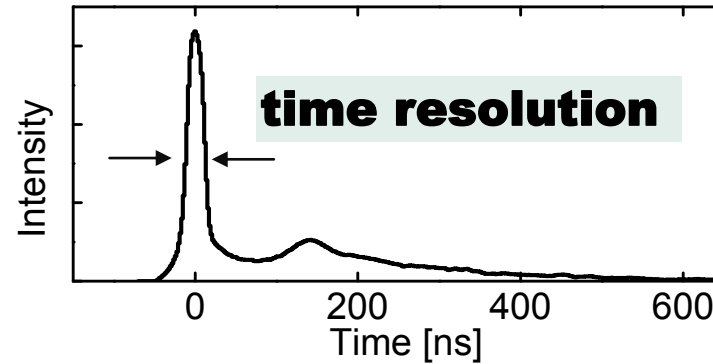
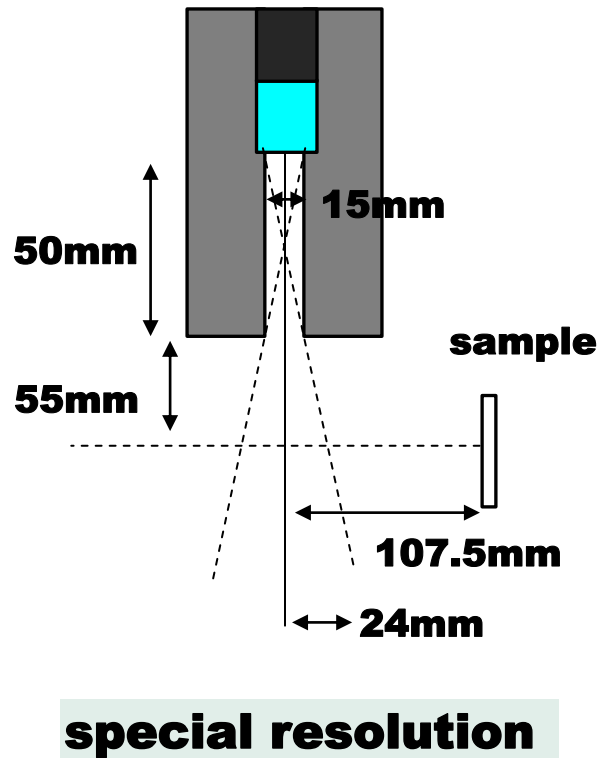
Calculated spin DOS



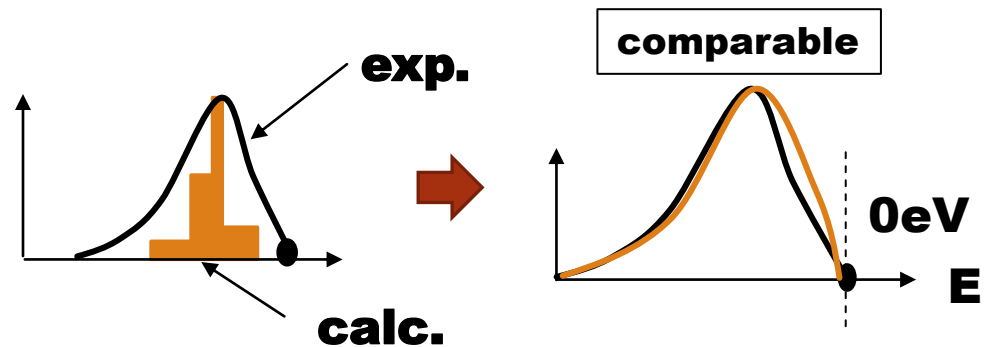
Experimental result

Convolution of energy resolution

- ✓ **Computational DOS had to be convolved at the resolution of the experimental setup.**

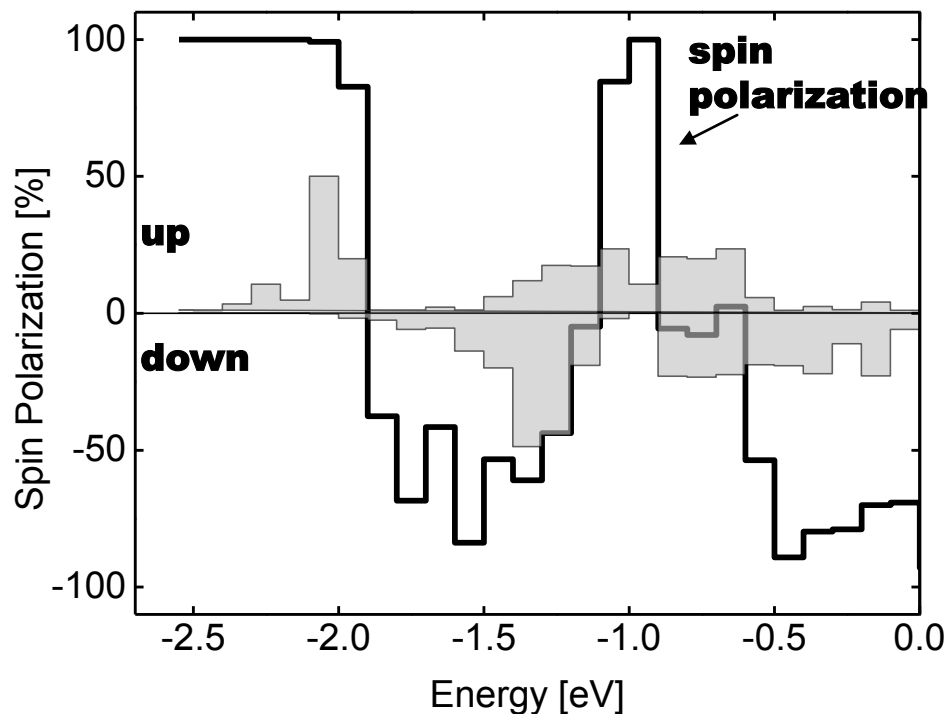


convolution of energy resolution

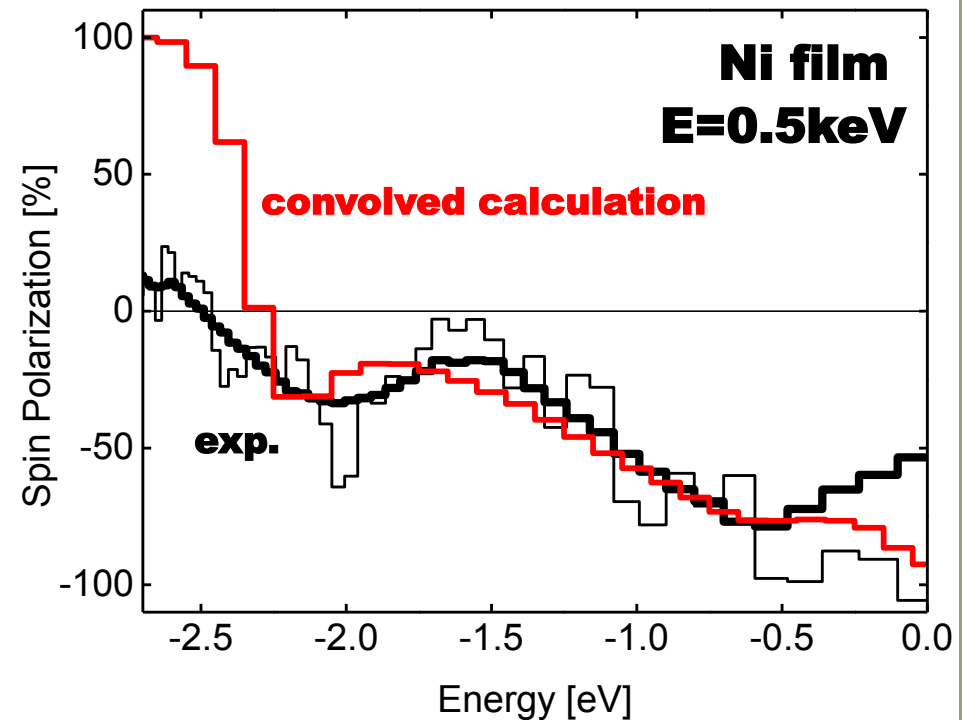


SP-PsTOF result (Ni film: magnetic)

- ✓ For the Ni thin film, a negative polarization was detected near the Fermi surface.
- ✓ well reproduced with experimental result



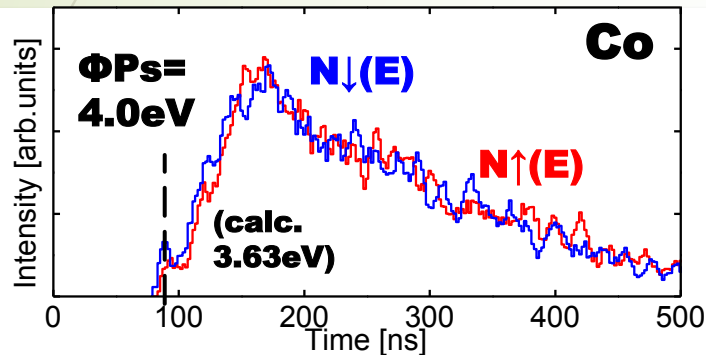
Calculated spin DOS



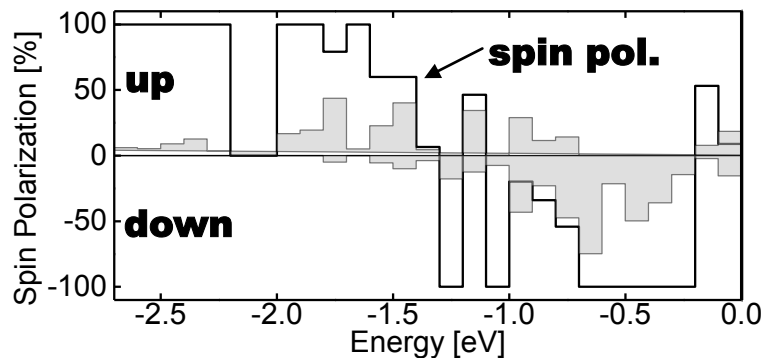
Experimental result

SP-PsTOF result (Co film: magnetic)

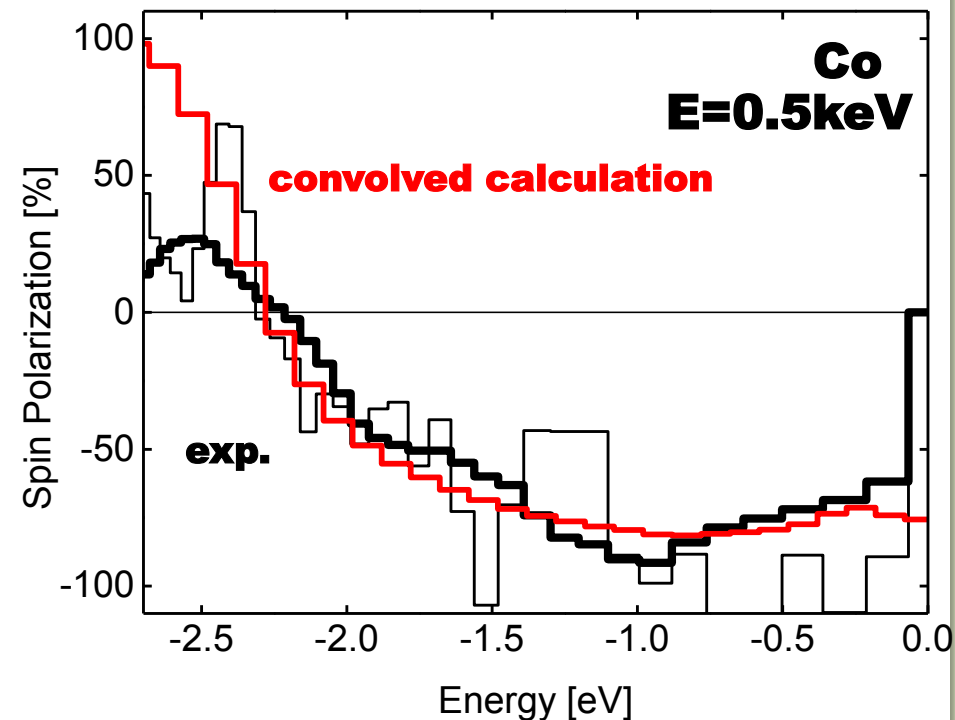
- ✓ As another example of a ferromagnetic material, Co film was also measured.
- ✓ Like Ni sample, a negative polarization was observed.



PsTOF time spectrum



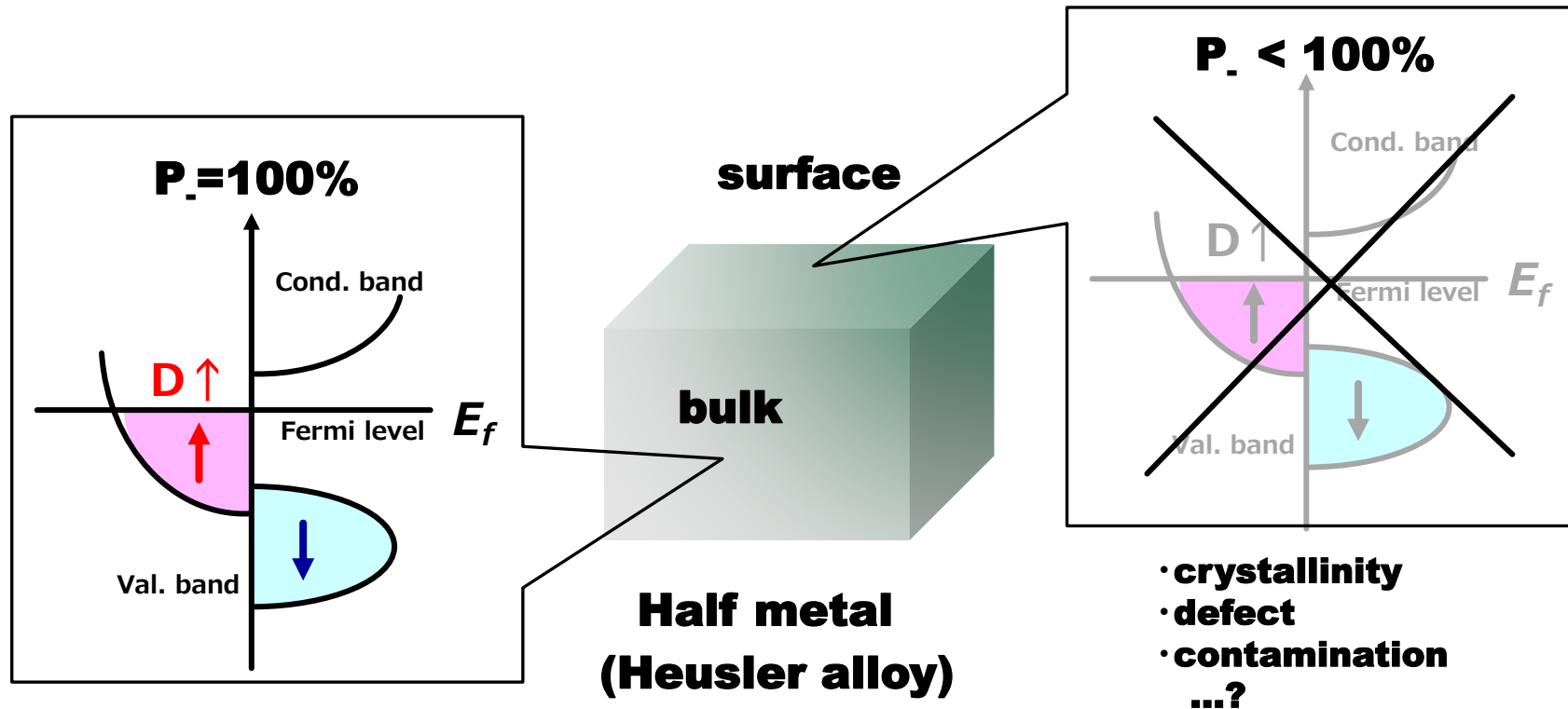
Calculated spin DOS



Energy-resolved topmost electron spin polarization

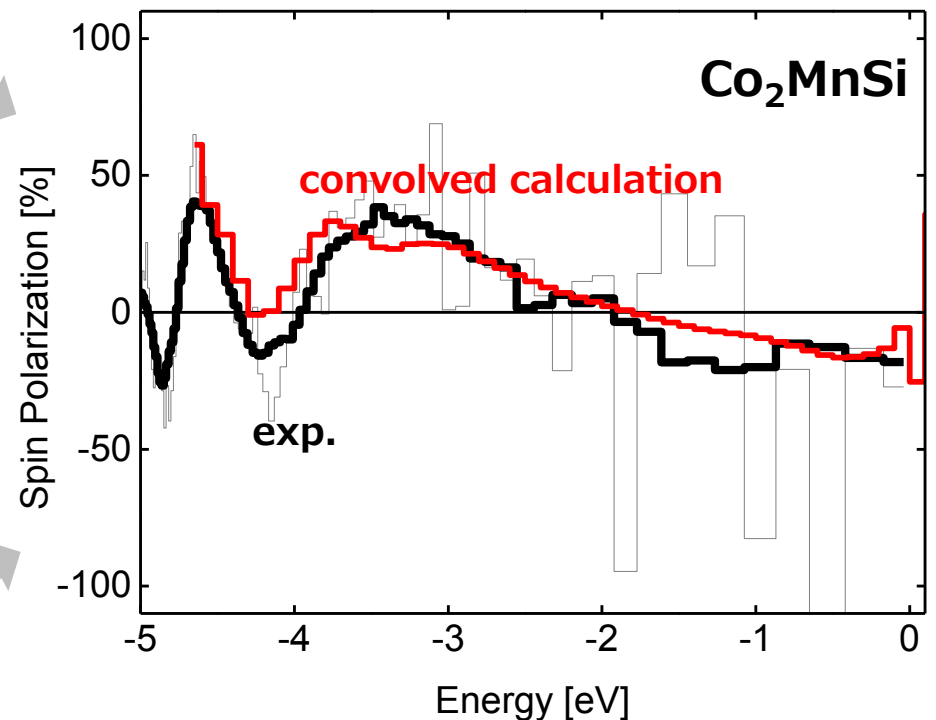
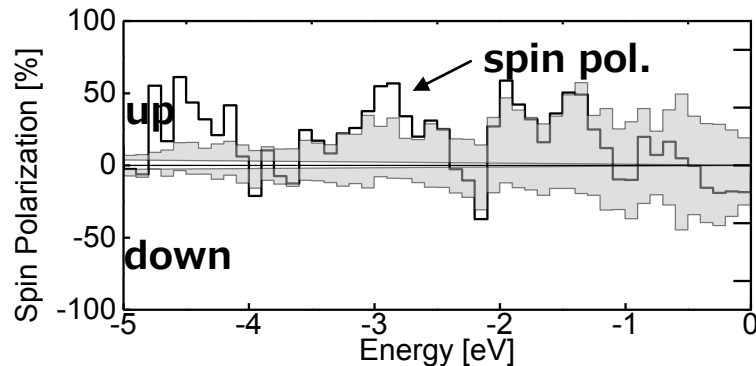
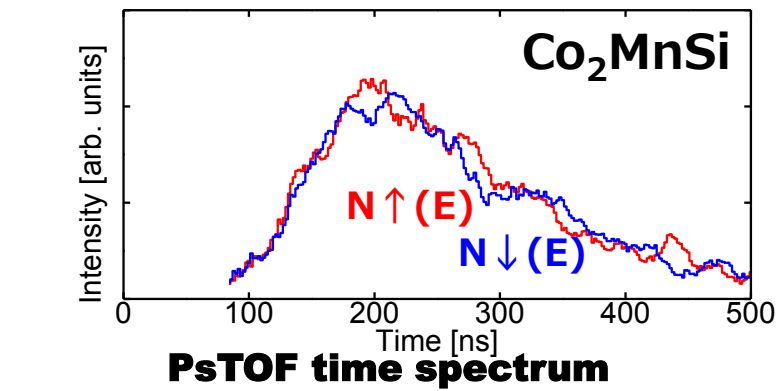
SP-PsTOF for Half metal

- ✓ **Half metal has a 100% polarization near the Fermi level in bulk, but not on the surface. The reason is still unclear.**
- ✓ **The SP-PsTOF measurement is suitable for evaluating the surface spin polarization of half metal.**



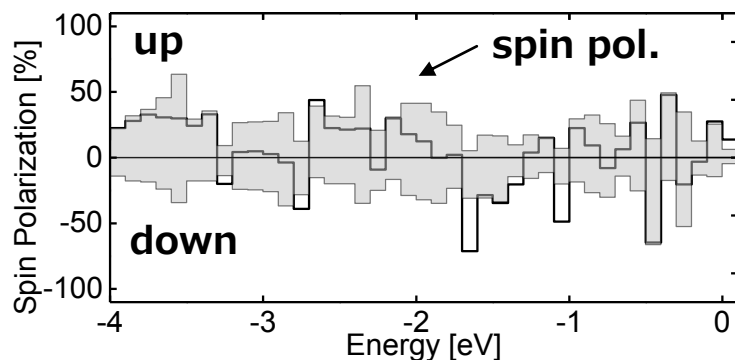
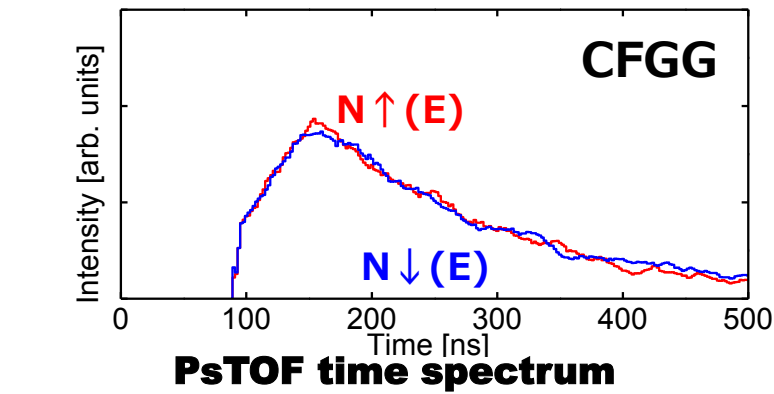
SP-PsTOF result (Co_2MnSi film)

- ✓ Energy-resolved spin polarization of Co_2MnSi (CMS) Heusler alloy film was obtained.
- ✓ Spin polarization near the Fermi level is very weak.

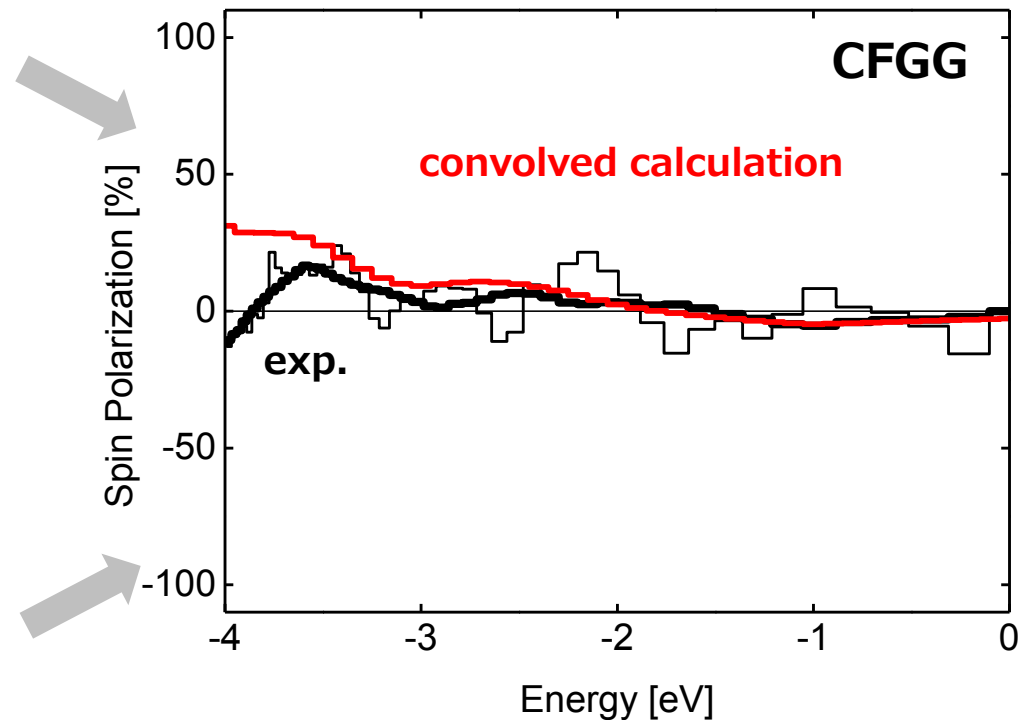


SP-PsTOF result ($\text{Co}_2\text{FeGa}_{0.5}\text{Ge}_{0.5}$ film)

- ✓ Energy-resolved spin polarization of $\text{Co}_2\text{FeGa}_{0.5}\text{Ge}_{0.5}$ (CFGG) Heusler alloy film was obtained.
- ✓ Spin polarization of surface electrons is very small.



Calculated spin DOS

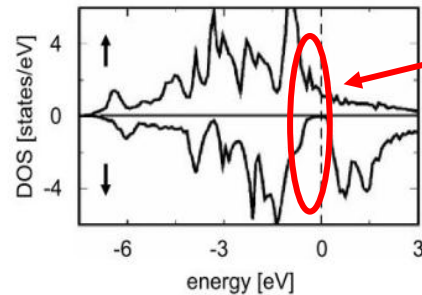


Energy-resolved topmost electron spin polarization

Improvement of SP-PsTOF apparatus

✓ For better measurements, improving the energy resolution is unavoidable.

Co₂MnSi spin-polarized DOS

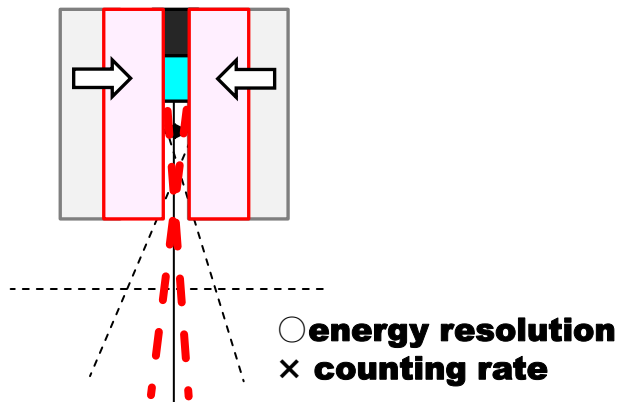


Below Fermi level $\sim 0.5\text{eV}$ is polarized
 → Energy resolution of $\sim 0.5\text{eV}$ is required.

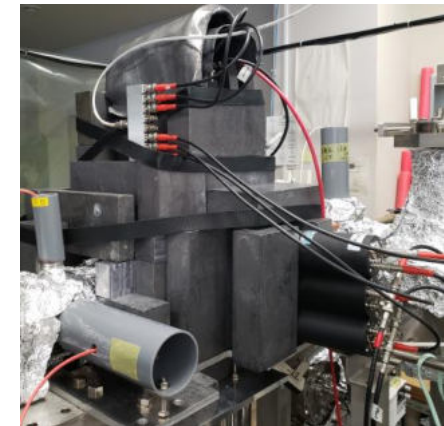
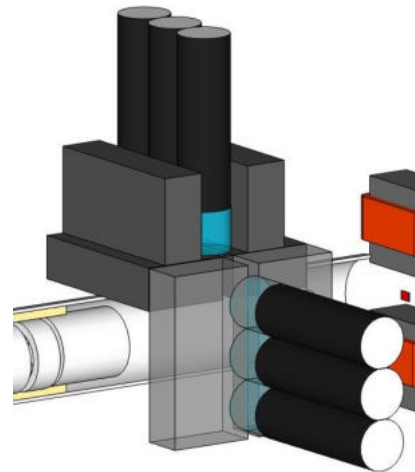


Present energy spread 1.4eV @ 3eV

Limitation of lead slit width



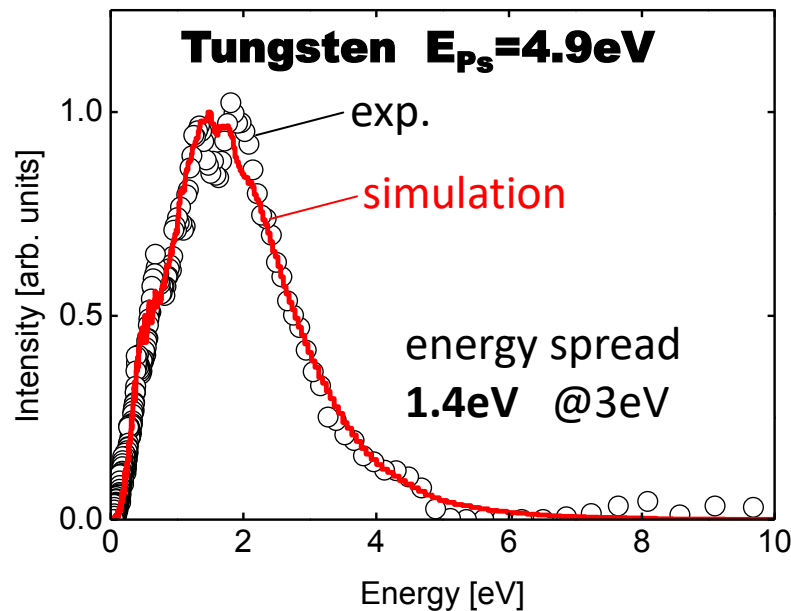
Additional detectors were installed



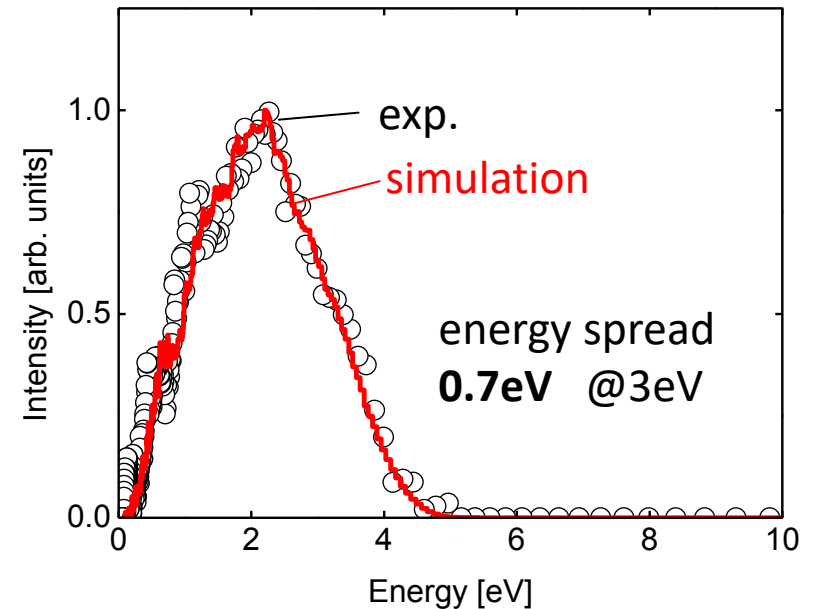
Improvement of PsTOF spectrum

✓ We try to find the best balance between counting rate and energy resolution.

slit width=15 mm
detector= 3 pcs.
Est. time resolution=30 ns
counting rate= 0.2-0.3 cps



slit width=8 mm
detector= 6 pcs.
Est. time resolution=5 ns
counting rate= 0.05-0.1 cps



Summary

- ✓ **Spin-polarized Ps TOF apparatus has been constructed.**
- ✓ **We succeeded in obtaining the first SP-PsTOF.**
 - **Clear asymmetry for ferromagnetic sample near $E=\Phi P_s$.**
- ✓ **In the Future, we have plan to measure the spintronics materials.**

Thank you for your attention!

Spin-Polarized Positronium Time-of-Flight Spectroscopy for Probing Spin-Polarized Surface Electronic StatesM. Maekawa¹, A. Miyashita¹, S. Sakai¹, S. Li¹, S. Entani¹, and A. Kawasuso^{1*}*National Institutes for Quantum and Radiological Science and Technology, 1233 Watanuki, Takasaki, Gunma 370-1292, Japan*Y. Sakuraba²*National Institute for Materials Science (NIMS), 1-2-1, Sengen, Tsukuba-city, Ibaraki 305-0047, Japan* (Received 5 November 2020; accepted 25 March 2021; published 4 May 2021)

The energy spectrum of positronium atoms generated at a solid surface reflects the electron density of states (DOS) associated solely with the first surface layer. Using spin-polarized positrons, the spin-dependent surface DOS can be studied. For this purpose, we have developed a spin-polarized positronium time-of-flight spectroscopy apparatus based on a ^{22}Na positron source and an electrostatic beam transportation system, which enables the sampling of topmost surface electrons around the Γ point and near the Fermi level. We applied this technique to nonmagnetic Pt(111) and W(001), ferromagnetic Ni(111), Co(0001) and graphene on them, $\text{Co}_2\text{FeGa}_{0.5}\text{Ge}_{0.5}$ (CFGG) and Co_2MnSi (CMS). The results showed that the electrons of Ni(111) and Co(0001) surfaces have characteristic negative spin polarizations, while these spin polarizations vanished upon graphene deposition, suggesting that the spin polarizations of graphene on Ni(111) and Co(0001) were mainly induced at the Dirac points that were out of range in the present measurement. The CFGG and CMS surfaces also exhibited only weak spin polarizations suggesting that the half-metallicity expected for these bulk states was not maintained at the surfaces.

DOI: [10.1103/PhysRevLett.126.186401](https://doi.org/10.1103/PhysRevLett.126.186401)

Slow positrons injected into the subsurface region of a metal diffuse back to the surface/vacuum interface and are emitted as positronium (Ps) atoms by picking up the outermost surface electrons when the Ps formation poten-

elucidate the nature of the spin polarization of the top-surface electronic states and play a valuable role in the field of spintronics.

Spin-polarized Ps spectroscopy was first demonstrated



Field asymmetry of o-Ps intensity

Number of **positrons** and **electrons** having up(\uparrow) and down(\downarrow) spins

$$n^{\uparrow} = \frac{1 + P_+}{2}$$

$$n^{\downarrow} = \frac{1 - P_+}{2}$$

$$n^{\uparrow} = D_{maj}(E)$$

$$n^{\downarrow} = D_{min}(E)$$

P_+ : Positron spin polarization

$D(E)$: electron density of state

Number of o-Ps

magnetic field	$ 11\rangle$	$ 1-1\rangle$	$ 10\rangle$
negative	$n^{\uparrow}n^{\uparrow}$	$n^{\downarrow}n^{\downarrow}$	$1/2(n^{\uparrow}n^{\downarrow} + n^{\downarrow}n^{\uparrow})$
positive	$n^{\uparrow}n^{\downarrow}$	$n^{\downarrow}n^{\uparrow}$	$1/2(n^{\uparrow}n^{\uparrow} + n^{\downarrow}n^{\downarrow})$

Fraction of o-Ps

$$F^{oPs} = \varepsilon_1(F_{|11\rangle} + F_{|1-1\rangle}) + \varepsilon_0 F_{|10\rangle}$$

($\varepsilon_1, \varepsilon_0$: detection efficiency)

Spectrum asymmetry

$$A^{oPs} = \frac{F_{Neg.field}^{oPs} - F_{Pos.field}^{oPs}}{F_{Neg.field}^{oPs} + F_{Pos.field}^{oPs}} = \frac{P_+(D_{maj}(E) - D_{min}(E)) \times \left(\varepsilon_1 - \frac{\varepsilon_0}{2}\right)}{(D_{maj}(E) + D_{min}(E)) \times \left(\varepsilon_1 + \frac{\varepsilon_0}{2}\right)}$$

$P_-(E)$: spin-polarized electron density of state

$$P_-(E) \propto A^{oPs}$$

$$= \frac{(2\varepsilon_1 - \varepsilon_0)}{(2\varepsilon_1 + \varepsilon_0)} P_+ \times P_-(E)$$

$$P_-(E) = \frac{D_{maj}(E) - D_{min}(E)}{D_{maj}(E) + D_{min}(E)}$$

actually,

$$P_-(E) \propto \frac{1}{f_{Ps}(E)} A^{oPs}(E)$$

f_{Ps} : Level dependence of Ps formation probability